

The product of the two eigenvalues is the constant  $\mu_+\mu_- = \det \mathcal{M} = -m^2$  so as  $\mu_-$  goes down,  $\mu_+$  must go up. **Minkowski, Yanagida, and Gell-Mann, Ramond, and Slansky** invented this “**seesaw**” mechanism as an explanation of why neutrinos have such small masses, less than 1 eV/ $c^2$ . If  $mc^2 = 10$  MeV, and  $\mu_-c^2 \approx 0.01$  eV, which is a plausible light-neutrino mass, then the rest energy of the huge mass would be  $Mc^2 = 10^7$  GeV. This huge mass would point at new physics, beyond the standard model. Yet the small masses of the neutrinos may be related to the weakness of their interactions.  $\square$

If we return to the orthogonal transformation (1.304) and multiply column  $\ell$  of the matrix  $O$  and row  $\ell$  of the matrix  $O^T$  by  $\sqrt{|R_\ell^{(d)}|}$ , then we arrive at the **congruency transformation** of Sylvester’s theorem

$$R = C \hat{R}^{(d)} C^T \quad (1.306)$$

in which the diagonal entries  $\hat{R}_\ell^{(d)}$  are either  $\pm 1$  or 0 because the matrices  $C_{k\ell} = \sqrt{|R_\ell^{(d)}|} O_{k\ell}$  and  $C^T$  have absorbed the **factors**  $|R_\ell^{(d)}|$ .

**Example 1.40** (Equivalence Principle) If  $G$  is a real, symmetric  $4 \times 4$  matrix then there’s a real  $4 \times 4$  matrix  $D = C^{T-1}$  such that

$$G_d = D^T G D = \begin{pmatrix} g_1 & 0 & 0 & 0 \\ 0 & g_2 & 0 & 0 \\ 0 & 0 & g_3 & 0 \\ 0 & 0 & 0 & g_4 \end{pmatrix} \quad (1.307)$$

in which the diagonal entries  $g_i$  are  $\pm 1$  or 0. Thus there’s a real  $4 \times 4$  matrix  $D$  that casts the real nonsingular symmetric metric  $g_{ik}$  of spacetime at any given point into the diagonal metric  $\eta_{j\ell}$  of flat spacetime by the congruence

$$g_d = D^T g D = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} = \eta. \quad (1.308)$$

Usually one needs different  $D$ s at different points. Since one can implement the congruence by changing coordinates, it follows that in any gravitational field, one may choose free-fall coordinates in which all physical laws take the same form as in special relativity without acceleration or gravitation at least over suitably small volumes of space-time (section 11.39).  $\square$