

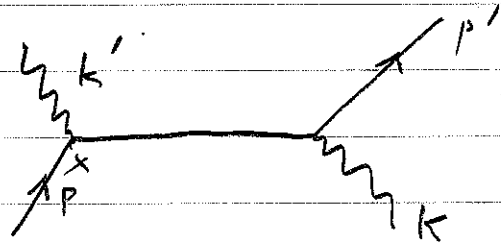
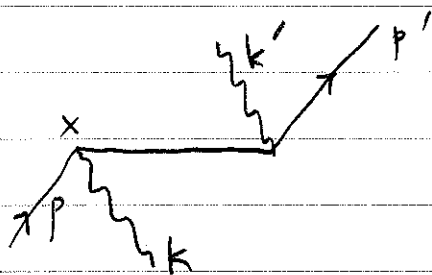
Compton Scattering with extra ± for fermions / bosons

$$-i \Delta_{\ell m}(x, y) = \langle 0 | \overline{\psi} \{ \psi_{\ell}(x) \psi_m^{\dagger}(y) \} | 0 \rangle$$

$$= -i \int \frac{d^4 q}{(2\pi)^4} \frac{P_{\ell m}^{(\ell)}(q) e^{iq(x-y)}}{q^2 + m^2 - i\epsilon}$$

$$P_{\ell m}^{(\ell)}(q) = P_{\ell m}(q) = [(-i \not{q} + m) \beta]_{\ell m}$$

$$\langle p' k' | S | p k \rangle = \frac{(-i)^2}{2} \langle p' k' | \int d^4 x d^4 y (ie)^2 (\overline{\psi} \not{A} \psi)_x (\overline{\psi} \not{A} \psi)_y | p k \rangle$$



take p absorbed at x

$$= e^2 \langle p' k' | \int d^4 x d^4 y (\overline{\psi} \not{A} \psi)_x (\overline{\psi} \not{A} \psi)_y | p k \rangle$$

$$\psi_{\alpha}(x) = \int \frac{d^3 p}{(2\pi)^{3/2}} \left[u_{\alpha}(p, s) a(p, s) e^{ipx} + v_{\alpha}(p, s) a^{\dagger}(p, s) e^{-ipx} \right]$$

$$A_{\mu}(x) = \int \frac{d^3 p}{(2\pi)^3 2p_0} \sum_s \left[e^{-ipx} e^{\mu}(p, s) a(p, s) + e^{ipx} e^{\mu}(p, s) a^{\dagger}(p, s) \right]$$

$$S = e^2 \langle 0 | a(p,s) a(k',\lambda') \int d^4x d^4y \theta(x-y) (\bar{\psi} \not{A} \psi)_x (\bar{\psi} \not{A} \psi)_y + \theta(y-x) (\bar{\psi} \not{A} \psi)_y (\bar{\psi} \not{A} \psi)_x a^\dagger(p,s) a^\dagger(k,\lambda) | 0 \rangle$$

$$= \frac{e^2 \langle 0 | a' a'}{(2\pi)^{3/2}} \int d^4x d^4y e^{i p x} \left[(\bar{\psi} \not{A})_x u(p,s) \theta(x-y) (\bar{\psi} \not{A} \psi)_y + \theta(y-x) (\bar{\psi} \not{A} \psi)_y (\bar{\psi} \not{A} u(p,s))_x \right] a^\dagger(k,\lambda) | 0 \rangle$$

$$= \frac{e^2 \langle 0 | a(k',\lambda')}{(2\pi)^3} \int d^4x d^4y e^{i p x - i p' y} \left[-\theta(x-y) (\bar{\psi} \not{A} u)_x \bar{u}' (\not{A} \psi)_y + \theta(y-x) \bar{u}' (\not{A} \psi)_y (\bar{\psi} \not{A})_x u \right] a^\dagger(k,\lambda) | 0 \rangle$$

$$= \frac{e^2 \langle 0 |}{(2\pi)^6 2 \sqrt{k^0 k'^0}} \int d^4x d^4y e^{i p x - i p' y} \left[-\theta(x-y) \bar{\psi}_x \not{\epsilon} u \bar{u}' \not{\epsilon} \psi_y + \theta(y-x) \bar{u}' \not{\epsilon} \psi_y \bar{\psi}_x \not{\epsilon} u \right] e^{i k x - i k' y} + \left[-\theta(x-y) \bar{\psi}_x \not{\epsilon}^* u \bar{u}' \not{\epsilon} \psi_y + \theta(y-x) \bar{u}' \not{\epsilon} \psi_y \bar{\psi}_x \not{\epsilon}^* u \right] e^{i k y - i k' x} | 0 \rangle$$

$$S = \frac{e^2}{(2\pi)^6 2\sqrt{k^0 k'^0}} \int d^4x d^4y \left[e^{ipx - ip'y + ikx - ik'y} \right. \\ \left. - i \Delta_{\alpha\beta}(\gamma-x) (\bar{u}' \not{\epsilon}'^*)_{\alpha} (\not{\epsilon} u)_{\beta} \right. \\ \left. + e^{ipx - ip'y - ikx + ik'y} \bar{u}' \not{\epsilon} (-i) \Delta(\gamma-x) \not{\epsilon}'^* u \right]$$

$$= \frac{-ie^2}{(2\pi)^6 2\sqrt{k^0 k'^0}} \int d^4x d^4y e^{ipx - ip'y} \\ \left[e^{ikx - ik'y} \bar{u}' \not{\epsilon}'^* \Delta(\gamma-x) \not{\epsilon} u \right. \\ \left. + e^{-ikx + ik'y} \bar{u}' \not{\epsilon} \Delta(\gamma-x) \not{\epsilon}'^* u \right]$$

$$= \frac{-ie^2}{(2\pi)^{10} 2\sqrt{k^0 k'^0}} \int d^4x d^4y d^4q \frac{e^{ipx - ip'y + iq(\gamma-x)}}{q^2 + m^2 - i\epsilon}$$

$$\left[e^{ikx - ik'y} \bar{u}' \not{\epsilon}'^* (-i \not{q} + m) \not{\epsilon} u \right. \\ \left. + e^{-ikx + ik'y} \bar{u}' \not{\epsilon} (-i \not{q} + m) \not{\epsilon}'^* u \right]$$

$$\begin{aligned}
S &= \frac{-ie^2}{(2\pi)^2 2\sqrt{k^0 k'^0}} \int d^4 q \left[\frac{\delta^4(p-q+k) \delta^4(-p'+q-k') \bar{u}' \not{k} (-i\not{q} + m) \not{k} u}{q^2 + m^2 - i\epsilon} \right. \\
&\quad \left. + \frac{\delta^4(p-q-k') \delta^4(-p'+q+k) \bar{u}' \not{k}' (-i\not{q} + m) \not{k} u}{q^2 + m^2 - i\epsilon} \right] \\
&= \frac{-ie^2 \delta^4(p+k-p'-k')}{(2\pi)^2 2\sqrt{k^0 k'^0}} \bar{u}' \left[\frac{\not{k}' (-i(\not{p} + \not{k}) + m) \not{k}}{(p+k)^2 + m^2 - i\epsilon} \right. \\
&\quad \left. + \frac{\not{k} (-i(\not{p} - \not{k}') + m) \not{k}'}{(p-k')^2 + m^2 - i\epsilon} \right] u
\end{aligned}$$

The denominators are, since $p^2 = -m^2$ and $k^2 = 0$,

$$(p+k)^2 + m^2 = -m^2 + 2p \cdot k + m^2 = 2p \cdot k$$

$$(p-k')^2 + m^2 = -m^2 - 2p \cdot k' + m^2 = -2p \cdot k'$$

In the lab. frame $p = (\vec{0}, m)$, and so these are

$$2p \cdot k = -2m\omega \quad \text{and} \quad -2p \cdot k' = 2m\omega'$$

Neither vanishes, and we drop the $-i\epsilon$'s.

B₉ (3.3.2)

$$S = -2\pi i \delta^4(p+k-p'-k') M \quad S_0$$

$$M = \frac{e^2}{4(2\pi)^3 \sqrt{k^0 k'^0}} \bar{u}(p's') \left\{ \not{\epsilon}' \frac{[-i(\not{p} + \not{k}) + m]}{p \cdot k} \not{\epsilon} - \frac{\not{\epsilon} [-i(\not{p} - \not{k}') + m]}{p \cdot k'} \right\} u(p's)$$

B₉ (3.4.15)

$$d\sigma = (2\pi)^4 u^{-1} |M|^2 \delta^4(p'+k-p-k') d^3p' d^3k'$$

where $u = \frac{1}{p \cdot k}$
 $p \cdot k'$

which in the lab. frame is

$$u = \frac{m\omega}{m\omega} = 1$$

So the S^3 cancels the d^3p' leaving
 $d\sigma = (2\pi)^4 |M|^2 \delta(\sqrt{(k-k')^2 + m^2} + \omega' - m - \omega) d^3k'$

The energy condition is

$$\omega^2 + \omega'^2 - 2\omega\omega' \cos\theta + m^2 = m + \omega - \omega'$$

one finds

$$w' \equiv w_c(\theta) = \frac{m}{m + w(1 - \cos\theta)}$$

Now

$$\begin{aligned} \int \delta(f(E) - X) df(E) &= \int \delta(E - E_x) dE = 1 \\ &= \int \delta(f(E) - X) \frac{df}{dE} dE \quad \text{So} \end{aligned}$$

$$\delta(f(E) - X) = \frac{\delta(E - E_x)}{|df/dE|}$$

Here $f(E) = p^{0'}(w') + w'$

$$\frac{d(p^{0'} + w')}{dw'} = \frac{d}{dw'} \left(\sqrt{w^2 + w'^2 - 2ww' \cos\theta + m^2} + w' \right)$$

$$= 1 + \frac{w' - w \cos\theta}{\sqrt{w^2 + w'^2 - 2ww' \cos\theta + m^2}} = \left| \frac{w' - w \cos\theta}{p^{0'}} + 1 \right|$$

So

$$\delta(p^{0'} + w' - m - w) = \frac{\delta(w' - w_c(\theta))}{\left| \frac{w' - w \cos\theta}{p^{0'}} + 1 \right|} = \frac{p^{0'} \delta(w' - w_c(\theta))}{|w' - w \cos\theta + p^{0'}|}$$

$$= \frac{p^{0'} \delta(w' - w_c(\theta))}{|m + w(1 - \cos\theta)|} = \frac{p^{0'} \delta(w' - w_c(\theta))}{mw/w'}$$

n

$$\delta(p^0 + w' - p^0 - w) = \frac{p^0 w'}{m w} \delta(w' - w c / \theta)$$

$$d^3 k' = w'^2 dw' d\Omega \quad \text{so}$$

$$d\sigma = (2\pi)^4 / |M|^2 \frac{p^0 w'^3}{m w} d\Omega$$

with

$$p^0 = m + w - w' \quad \text{and} \quad w' = w c / \theta.$$

usually, we average over initial spins & polarizations and sum over final ones.

$$M = \bar{u}'_\alpha A_{\alpha\beta} u_\beta$$

$$|M|^2 = M M^* = \bar{u}'_\alpha A_{\alpha\beta} u_\beta u'_r (A^\dagger_\beta)_rs u'_s$$

$$= \bar{u}'_\alpha A_{\alpha\beta} u_\beta \bar{u}'_r (\beta A^\dagger_\beta)_rs u'_s$$

$$= A_{\alpha\beta} u_\beta \bar{u}'_r (\beta A^\dagger_\beta)_rs u'_s \bar{u}'_\alpha$$

$$= \text{Tr} A u \bar{u} \beta A^\dagger \beta u \bar{u}$$

$$\sum_s u_\beta(p,s) \bar{u}'_r(p,s) = \frac{(-i\not{p} + m)_{\beta\gamma}}{2p^0}$$

$$\sum_{s'} u_s(p's') \bar{u}_\alpha(p's') = \frac{(-i \not{p} + m) \delta_\alpha}{2p^0}$$

Our γ 's are

$$\gamma^0 = -i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \vec{\gamma} = -i \begin{pmatrix} 0 & \vec{\sigma} \\ -\vec{\sigma} & 0 \end{pmatrix} \quad (5.4.17)$$

So $\gamma^{0\dagger} = -\gamma^0 \quad \vec{\gamma}^\dagger = \vec{\gamma}$

$$\{\gamma^\mu, \gamma^\nu\} = 2\eta^{\mu\nu}$$

so

$$\beta \gamma^{\mu\dagger} \beta = -\gamma^\mu$$

Now

$$A = \not{\epsilon}^{x'} [-i(\not{p} + \not{k}) + m] \not{\epsilon} / p \cdot k$$

$$- \not{\epsilon} [-i(\not{p} - \not{k}) + m] \not{\epsilon}^{x'} / p \cdot k'$$

and

$$A \beta \not{\epsilon} \beta = \beta p_\mu \gamma^{\mu\dagger} \beta = -\not{p}$$

$$\beta \not{\epsilon} \beta = \beta \epsilon_m^* \gamma^{\mu\dagger} \beta = -\not{\epsilon}^*$$

$$\beta (\not{\epsilon}^*)^\dagger \beta = \beta \epsilon_\mu \gamma^{\mu\dagger} \beta = -\not{\epsilon}$$

So

$$A^\dagger = \not{\epsilon}^\dagger [+i(\not{p}^\dagger + \not{k}^\dagger) + m] (\not{\epsilon}^\dagger)^{\dagger} / p \cdot k$$

$$- (\not{\epsilon}^{\dagger})^\dagger [i(\not{p}^\dagger - \not{k}^\dagger) + m] (\not{\epsilon}^\dagger)^\dagger / p \cdot k'$$

So

$$\beta A^\dagger \beta = \beta \not{\epsilon}^\dagger \beta \beta [i(\not{p}^\dagger + \not{k}^\dagger) + m] \beta \beta (\not{\epsilon}^{\dagger})^\dagger \beta / p \cdot k$$

$$- \beta (\not{\epsilon}^{\dagger})^\dagger \beta \beta [i(\not{p}^\dagger - \not{k}^\dagger) + m] \beta \beta \not{\epsilon}^\dagger / p \cdot k'$$

$$= - \not{\epsilon}^* [-i(\not{p} + \not{k}) + m] (-\not{\epsilon}') / p \cdot k$$

$$+ \not{\epsilon}' [-i(\not{p} - \not{k}') + m] (-\not{\epsilon}^*) / p \cdot k'$$

So

$$\sum_{ss'} |M|^2 = \frac{e^4}{64(2\pi)^6 \omega \omega' p^0 p'^0}$$

$$\times \text{Tr} \left\{ \frac{\not{\epsilon}^{\dagger} [-i(\not{p} + \not{k}) + m] \not{\epsilon}}{p \cdot k} - \frac{\not{\epsilon} [-i(\not{p} - \not{k}') + m] \not{\epsilon}^{\dagger}}{p \cdot k'} \right\}$$

$$(-i\not{p} + m) \left\{ \frac{\not{\epsilon}^* [-i(\not{p} + \not{k}) + m] \not{\epsilon}'}{p \cdot k} - \frac{\not{\epsilon}' [-i(\not{p} - \not{k}') + m] \not{\epsilon}^{\dagger}}{p \cdot k'} \right\}$$

$$(-i\not{k}' + m)$$

In Coulomb's gauge in lab. frame

$$e \cdot p = e^* \cdot p = e' \cdot p = e'^* \cdot p = 0$$

since $e^0 = e'^0 = 0$ & $\vec{p} = 0$.

$$\begin{aligned} Xg &= f_{\mu\nu} g_{\rho\sigma} \gamma^{\mu} \gamma^{\nu} \\ &= f_{\mu\nu} g_{\rho\sigma} (-\gamma^{\nu} \gamma^{\mu} + 2\eta^{\mu\nu}) \\ &= -g\cancel{f} + 2fg \quad (X^2 = f \cdot f) \quad \text{so} \end{aligned}$$

$$\begin{aligned} (-i\not{p} + m)\cancel{\phi}(-i\not{p} + m) &= \cancel{\phi}(i\not{p} + m)(-i\not{p} + m) \\ &= \cancel{\phi}(\not{p}^2 + m^2) = \cancel{\phi}(-m^2 + m^2) = 0. \end{aligned}$$

This works also for e' , e^* , & e'^* .

One gets

$$\sum_{ss'} |M|^2 = \frac{e^4}{64(2\pi)^6 \omega \omega' p^0 p'^0}$$

$$\begin{aligned} & \times \text{Tr} \left[\frac{\cancel{\epsilon}^* \not{k} \cancel{\phi}}{p \cdot k} + \frac{\cancel{\phi} \not{k}' \cancel{\epsilon}'^*}{p \cdot k'} \right] (-i\not{p} + m) \\ & \times \left[\frac{\cancel{\phi}' \not{k} \cancel{\phi}}{p \cdot k} + \frac{\cancel{\epsilon}' \not{k}' \cancel{\phi}^*}{p \cdot k'} \right] (-i\not{p}' + m) \end{aligned} \quad (8.7.23)$$

Now

$$\text{Tr}(\gamma_{\mu_1} \gamma_{\mu_2} \dots \gamma_{\mu_{2m+1}}) = 0$$

because

$$\gamma_5 \gamma_{\mu} \gamma_5^{-1} = -\gamma_{\mu}$$

where $\gamma_5 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \gamma_5^{-1}$

Thus

$$\text{Tr} \gamma_{\mu_1} \gamma_{\mu_2} \dots \gamma_{\mu_{2m+1}}$$

$$= \text{Tr} \gamma_5 \gamma_{\mu_1} \gamma_5 \gamma_5 \gamma_{\mu_2} \gamma_5^2 \gamma_{\mu_3} \dots \gamma_5 \gamma_{\mu_{2m+1}} \gamma_5$$

$$= (-1)^{2m+1} \text{Tr} \gamma_{\mu_1} \gamma_{\mu_2} \dots \gamma_{\mu_{2m+1}}$$

The appendix to ch 8 shows that

$$2 \text{Tr} \gamma_{\mu} \gamma_{\nu} = 2 \eta_{\mu\nu} \text{Tr} \mathbf{1} = 8 \eta_{\mu\nu}$$

$$\text{Tr} \gamma_{\mu} \gamma_{\nu} = 4 \eta_{\mu\nu}$$

$$\text{Tr} \gamma_{\mu} \gamma_{\nu} \gamma_{\rho} \gamma_{\sigma} = 4 (\eta_{\mu\nu} \eta_{\rho\sigma} - \eta_{\mu\rho} \eta_{\nu\sigma} + \eta_{\mu\sigma} \eta_{\nu\rho})$$

etc

$$\text{Also } e \cdot p = 0 \quad k \cdot h = k' \cdot h' = 0 \quad e \cdot e^\wedge = e' \cdot e'^\wedge = 1$$

We will use real e 's to further simplify the results. We then drop all x 's.

$$\not{x} \not{x} \not{x} = -\not{x} \not{x} \not{x} = -\not{x}$$

$$k \not{x} k = -k k \not{x} + k 2p \cdot h = 2k p \cdot h$$

So

$$T_1 \equiv \text{Tr } \not{e}' k \not{x} \not{x} \not{x} k \not{e}' k'$$

$$= -\text{Tr } \not{e}' k \not{x} k \not{e}' k'$$

$$= -2 p \cdot k \text{Tr } \not{e}' k \not{e}' k'$$

$$= -8 p \cdot h \left(2 e' \cdot h e' \cdot p' - h \cdot p' \right)$$

$$e' \cdot p' = e' \cdot (p + h - h') = e' \cdot h$$

$$k \cdot p' = -\frac{1}{2} (p' - h)^2 - \frac{1}{2} m^2 = -\frac{1}{2} (p - h')^2 - \frac{1}{2} m^2 = p \cdot k'$$

So

$$T_1 = -16 p \cdot h (e' \cdot h)^2 + 8 p \cdot h p \cdot h'$$

Eventually, we have

$$\sum_{ss'} |M|^2 = \frac{e^4}{64(2\pi)^6 \omega \omega' p^0 p'^0} \left[\frac{8(k \cdot k')^2}{k \cdot p \ k' \cdot p} + 32(e \cdot e')^2 \right]$$

in lab. frame

$$k \cdot k' = \omega \omega' (\cos \theta - 1) = m \omega \omega' \left(\frac{1}{\omega} - \frac{1}{\omega'} \right)$$

$$p \cdot k = -m \omega$$

$$p \cdot k' = -m \omega'$$

so

$$\frac{1}{2} \sum_{ss'} d\sigma = \frac{e^4 \omega'^2 d\Omega}{64 \pi^2 m^2 \omega^2} \left[\frac{\omega}{\omega'} + \frac{\omega'}{\omega} - 2 + 4(e \cdot e')^2 \right]$$

Klein & Nishina 1929

Average over initial polarization

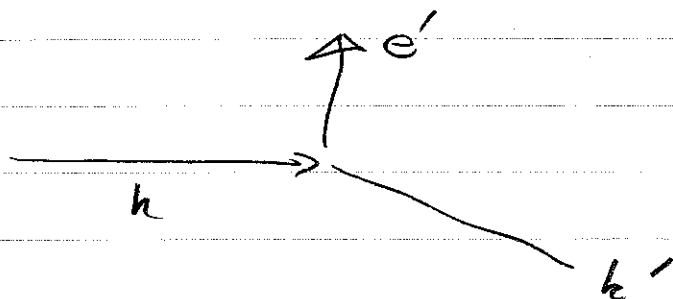
$$\frac{1}{2} \sum_{\lambda} e_i e_j = \frac{1}{2} (\delta_{ij} - \hat{k}_i \hat{k}_j)$$

$$\frac{1}{4} \sum_{ss'} d\sigma = \frac{e^4 \omega'^2 d\Omega}{64 \pi^2 m^2 \omega^2} \left(\frac{\omega}{\omega'} + \frac{\omega'}{\omega} - 2(\hat{k} \cdot \hat{e}')^2 \right)$$

which is highest when $\hat{k} \cdot \hat{e}' = 0$,

that is, $e' \cdot k = e' \cdot k' = 0$.

So e' tends to be \perp to plane of scattering



Usually, one sums over the e' 's:

$$\frac{1}{4} \sum_{SS' \lambda \lambda'} d\sigma = \frac{e^4 \omega'^2 dV \Omega}{32 \pi^2 m^2 \omega^2} \left(\frac{\omega}{\omega'} + \frac{\omega'}{\omega} - 1 + \cos^2 \theta \right)$$

where $\hat{k} \cdot \hat{k}' = \cos \theta$,

In non-rel case, $\omega \ll m$, then

$\omega \approx \omega'$ and

$$\frac{1}{4} \sum d\sigma = \frac{e^4 dV \Omega}{32 \pi^2 m^2} (1 + \cos^2 \theta)$$

$$\frac{1}{4} \sum \int d\Omega = \sigma = \frac{16\pi}{3} \frac{e^4}{32 \pi^2 m^2}$$

$$= \frac{e^4}{6\pi m^2} = \sigma_T \quad \text{Thomson}$$