

The interaction hamiltonian density is

$$\mathcal{H}_I = e \bar{\Psi} \gamma^\mu \frac{1}{2} \vec{\sigma} \cdot \vec{A}_\mu \Psi.$$

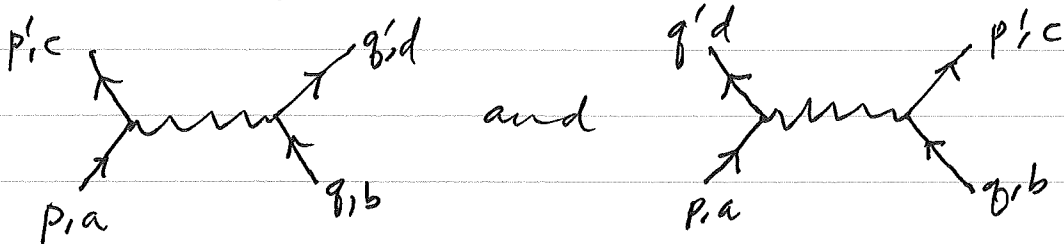
The propagator changes

$$-\frac{i\eta_{\mu\nu}}{q^2 + i\epsilon} \rightarrow -\frac{i\eta_{\mu\nu} \delta_{ij}}{q^2 + i\epsilon}$$

and the vertex changes

$$\begin{array}{c} \nearrow \\ \text{---} \\ \searrow \end{array} = -ie\gamma^\mu \rightarrow \begin{array}{c} \nearrow^b \\ \text{---} \\ \searrow^a \end{array} = -ie\gamma^\mu \frac{1}{2} \sigma_{ba}^i.$$

So the diagrams for  $f_a + f_b \rightarrow f_c + f_d$  are



$$i\mathcal{M} = \bar{u}(p'_c) (-ie\gamma^\mu \frac{1}{2} \sigma_{ca}^i) u(p_s) \frac{(-i\eta_{\mu\nu} \delta_{ij})}{(p \cdot p')^2 + i\epsilon} \bar{u}(q'_d) (-ie\gamma^\nu \frac{1}{2} \sigma_{db}^j) u(q, \epsilon)$$

$$- \bar{u}(q'_d) (-ie\gamma^\mu \frac{1}{2} \sigma_{da}^j) u(p_s) \frac{(-i\eta_{\mu\nu} \delta_{ij})}{(p - q')^2 + i\epsilon} \bar{u}(p'_c) (-ie\gamma^\nu \frac{1}{2} \sigma_{cb}^i) u(q, \epsilon)$$

$$i\mathcal{M} = \frac{ie^2}{4} \left[ \frac{\bar{u}(p',s') \gamma^\mu u(p,s) \sigma_{c\lambda}^i \sigma_{db}^i \bar{u}(q',t') \gamma_\mu u(q,t)}{(p-p')^2 + i\epsilon} - \frac{\bar{u}(q',t') \gamma^\mu u(p,s) \sigma_{d\lambda}^i \sigma_{cb}^i \bar{u}(p',s') \gamma_\mu u(q,t)}{(p-q')^2 + i\epsilon} \right].$$

For instance, if  $a=b=1$  and  $c=d=2$ , then

$$\sigma_{21}^i \sigma_{21}^i = 1^2 + i^2 = 1 + (-1) = 0,$$

and the amplitude vanishes.

Another example:  $a=1, b=2, c=1, d=2$ :

$$\sigma_{11}^i \sigma_{22}^i = \sigma_{11}^3 \sigma_{22}^3 = -1$$

$$\sigma_{21}^i \sigma_{12}^i = \sigma_{21}^1 \sigma_{12}^1 + \sigma_{21}^2 \sigma_{12}^2 = 2.$$