

Isospin: $[I_i, I_j] = i \sum_{k \in \{j, k\}} I_k$.

Isospin is dimensionless: no \hbar in $[I_i, I_j]$.
 Apart from the lack of \hbar , the algebra of isospin is the same as that of angular momentum.

The proton p and the neutron n form a doublet

$$\begin{pmatrix} p \\ n \end{pmatrix} \text{ with } |p\rangle = \left| \frac{1}{2}, \frac{1}{2} \right\rangle \text{ and } |n\rangle = \left| \frac{1}{2}, -\frac{1}{2} \right\rangle.$$

The π mesons form a triplet

$$|\pi^+\rangle = |1, 1\rangle \quad |\pi^0\rangle = |1, 0\rangle \quad |\pi^-\rangle = |1, -1\rangle.$$

The Λ hyperon is a singlet $|\Lambda\rangle = |0, 0\rangle$.

The 1232 MeV resonance in $\pi^+ p$ scattering

(Farrar, 1951) is a quartet

$$\begin{aligned} |\Delta^{++}\rangle &= \left| \frac{3}{2}, \frac{3}{2} \right\rangle & |\Delta^+\rangle &= \left| \frac{3}{2}, \frac{1}{2} \right\rangle \\ |\Delta^0\rangle &= \left| \frac{3}{2}, -\frac{1}{2} \right\rangle & |\Delta^-\rangle &= \left| \frac{3}{2}, -\frac{3}{2} \right\rangle. \end{aligned}$$

The up and down quarks form a doublet

$$|u\rangle = \left| \frac{1}{2}, \frac{1}{2} \right\rangle \quad |d\rangle = \left| \frac{1}{2}, -\frac{1}{2} \right\rangle.$$

Dynamics: The np states form a triplet

$$\begin{aligned} |1, 1\rangle &= |pp\rangle & |1, 0\rangle &= \frac{1}{\sqrt{2}} (|pn\rangle + |np\rangle) \\ |1, -1\rangle &= |nn\rangle \end{aligned}$$

and one singlet $|0, 0\rangle = \frac{1}{\sqrt{2}} (|pn\rangle - |np\rangle)$.

This last is the deuteron d , which is a stable bound state. The n - p interaction is something like

$$V = (\vec{I}_1 + \vec{I}_2)^2$$

Thus $\langle 1, m | V | 1, m \rangle = 2$ while $\langle 0, 0 | V | 0, 0 \rangle = 0$.

Consider $p+p \rightarrow d+\pi^+$

$p+n \rightarrow d+\pi^0$

$$|pp\rangle = |11\rangle \quad |d\rangle = |0,0\rangle \quad |\pi^+\rangle = |1,1\rangle$$

Now

$$|pn\rangle = \left| \frac{1}{2} \right\rangle \left| -\frac{1}{2} \right\rangle = \frac{1}{2} \left(\left| \frac{1}{2} \right\rangle \left| \frac{1}{2} \right\rangle + \left| \frac{1}{2} \right\rangle \left| -\frac{1}{2} \right\rangle - \left(\left| \frac{1}{2} \right\rangle \left| \frac{1}{2} \right\rangle - \left| \frac{1}{2} \right\rangle \left| -\frac{1}{2} \right\rangle \right)$$

$$= \frac{1}{\sqrt{2}} (|110\rangle - |100\rangle)$$

So since $[S, \vec{I}] = 0$

$$\langle d, \pi^+ | S | pp \rangle = \langle 11 | S | 11 \rangle \quad \text{w/wk}$$

$$\langle d, \pi^0 | S | pn \rangle = \langle 10 | S | \frac{1}{\sqrt{2}} (|110\rangle - |100\rangle) \rangle$$

$$= \frac{1}{\sqrt{2}} \langle 10 | S | 110 \rangle$$

So we expect the ratio of cross-sections to be

$$\frac{\sigma(pp \rightarrow d\pi^+)}{\sigma(pn \rightarrow d\pi^0)} = \frac{2}{2} \quad \text{as observed.}$$