

1. Consider an electron with orbital angular momentum $\ell = 1$ (and of course spin $1/2$) in a state in which the total angular momentum $J = 3/2$ and $M = 1/2$. (a) What is the probability that the spin of the electron is up? (b) Same question, but if $J = 1/2$?

$$(a) \quad |l + \frac{1}{2}, l - \frac{1}{2}\rangle = \frac{1}{\sqrt{2\ell + 1}} \left(\sqrt{2\ell} |l, l-1\rangle | \frac{1}{2}, + \rangle + |l, l\rangle | \frac{1}{2}, - \rangle \right)$$

Set $\ell = 1$

$$| \frac{3}{2}, \frac{1}{2} \rangle = \frac{1}{\sqrt{3}} \left(\sqrt{2} |1, 0\rangle | \frac{1}{2}, + \rangle + |1, 1\rangle | \frac{1}{2}, - \rangle \right)$$

So the probability that the electron's spin is up is

$$\left(\sqrt{\frac{2}{3}} \right)^2 = \frac{2}{3}$$

$$(b) \quad |l - \frac{1}{2}, l - \frac{1}{2}\rangle = \frac{1}{\sqrt{2\ell + 1}} \left(\sqrt{2\ell} |l, \ell\rangle | \frac{1}{2}, - \rangle - |l, \ell-1\rangle | \frac{1}{2}, + \rangle \right)$$

Set $\ell = 1$

$$| \frac{1}{2}, \frac{1}{2} \rangle = \frac{1}{\sqrt{3}} \left(\sqrt{2} |1, 1\rangle | \frac{1}{2}, - \rangle - |1, 0\rangle | \frac{1}{2}, + \rangle \right)$$

Now the spin-up probability is

$$\left(-\frac{1}{\sqrt{3}} \right)^2 = \frac{1}{3}$$

2. The Δ resonance occurs in pion-nucleon scattering at about 1232 MeV. The pion has isospin 1, so that $|\pi^+\rangle = |1, 1\rangle$, $|\pi^0\rangle = |1, 0\rangle$, and $|\pi^-\rangle = |1, -1\rangle$. The proton-neutron system has isospin 1/2 with $|p\rangle = |1/2, 1/2\rangle$ and $|n\rangle = |1/2, 1/2\rangle$. This giant resonance has $J = 3/2$ and isospin $I = 3/2$. The $I_z = 1/2$ state of this resonance can decay into both $|\pi^0, p\rangle$ and $|\pi^+, n\rangle$. What is the ratio of these two decay modes?

$$|3/2, 1/2\rangle = \frac{1}{\sqrt{3}} \left(\sqrt{2} |1, 0\rangle |1/2, +\rangle + |1, 1\rangle |1/2, -\rangle \right)$$

$$|\Delta^+\rangle = \frac{1}{\sqrt{3}} \left(\sqrt{2} |\pi^0\rangle |p\rangle + |\pi^+\rangle |n\rangle \right).$$

So

$$\langle \pi^0 p | \Delta^+ \rangle = \frac{\sqrt{2}}{\sqrt{3}}$$

and

$$\langle \pi^+ n | \Delta^+ \rangle = \frac{1}{\sqrt{3}}.$$

So

$$\frac{|\langle \pi^0 p | \Delta^+ \rangle|^2}{|\langle \pi^+ n | \Delta^+ \rangle|^2} = 2.$$

3. (a) Compute the derivatives of the unit vectors \hat{r} , $\hat{\theta}$, and $\hat{\phi}$ with respect to the variables r , θ , and ϕ . Your formulas should express these derivatives in terms of the basis vectors \hat{r} , $\hat{\theta}$, and $\hat{\phi}$. (b) Using the formulas of (a), derive a formula for the laplacian $\nabla \cdot \nabla$.

(a) First, we compute the derivatives of \hat{r} , $\hat{\theta}$, and $\hat{\phi}$ with respect to r , θ , and ϕ .

$$\frac{\partial \hat{r}}{\partial r} = 0$$

$$\frac{\partial \hat{r}}{\partial \theta} = \hat{x} \cos \theta \cos \phi + \hat{y} \cos \theta \sin \phi - \hat{z} \sin \theta = \hat{\theta}$$

$$\frac{\partial \hat{r}}{\partial \phi} = -\hat{x} \sin \theta \sin \phi + \hat{y} \sin \theta \cos \phi = \sin \theta \hat{\phi}$$

$$\frac{\partial \hat{\theta}}{\partial r} = 0$$

$$\frac{\partial \hat{\theta}}{\partial \theta} = -\hat{x} \sin \theta \cos \phi - \hat{y} \sin \theta \sin \phi - \hat{z} \cos \theta = -\hat{r}$$

$$\frac{\partial \hat{\theta}}{\partial \phi} = -\hat{x} \cos \theta \sin \phi + \hat{y} \cos \theta \cos \phi = \cos \theta \hat{\phi}$$

$$\frac{\partial \hat{\phi}}{\partial r} = 0$$

$$\frac{\partial \hat{\phi}}{\partial \theta} = 0$$

$$\frac{\partial \hat{\phi}}{\partial \phi} = -\hat{x} \cos \phi - \hat{y} \sin \phi = -\sin \theta \hat{r} - \cos \theta \hat{\theta}$$

(b) Now the Laplacian Δ is

$$\nabla \cdot \nabla = \left(\hat{r} \frac{\partial}{\partial r} + \frac{\hat{\theta}}{r} \frac{\partial}{\partial \theta} + \frac{\hat{\phi}}{r \sin \theta} \frac{\partial}{\partial \phi} \right) \cdot \left(\hat{r} \frac{\partial}{\partial r} + \frac{\hat{\theta}}{r} \frac{\partial}{\partial \theta} + \frac{\hat{\phi}}{r \sin \theta} \frac{\partial}{\partial \phi} \right)$$

$$= \frac{\partial^2}{\partial r^2} + \frac{\hat{\theta}}{r} \cdot \frac{\partial \hat{r}}{\partial \theta} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} + \frac{\hat{\phi}}{r \sin \theta} \cdot \frac{\partial \hat{r}}{\partial \phi} \frac{\partial}{\partial r}$$

$$+ \frac{\hat{\phi}}{r \sin \theta} \cdot \frac{\partial \hat{\theta}}{\partial \phi} \frac{\partial}{\partial \theta} + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \phi^2}$$

$$= \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{\cot \theta}{r} \frac{\partial}{\partial \theta} + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \phi^2}$$

$$= \frac{\partial^2}{\partial r^2} + \frac{2}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \left[\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right]$$

$$= \frac{1}{r^2} \left[\frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right]$$

4. Consider the two-body hamiltonian

$$H = \frac{p_1^2}{2m_1} + \frac{p_2^2}{2m_2} + V(r_1 - r_2). \quad (1)$$

Define the total momentum as

$$P = p_1 + p_2 \quad (2)$$

the relative momentum as

$$p = \frac{m_2 p_1 - m_1 p_2}{m_1 + m_2}. \quad (3)$$

and the separation as

$$r = r_1 - r_2. \quad (4)$$

Now show that H can be written as

$$H = \frac{P^2}{2M} + \frac{p^2}{2\mu} + V(r) \quad (5)$$

in which $M = m_1 + m_2$ is the total mass, and

$$\mu = \frac{m_1 m_2}{m_1 + m_2} \quad (6)$$

is the reduced mass.

$$\frac{P^2}{2M} + \frac{p^2}{2\mu} = \frac{(p_1 + p_2)^2}{2(m_1 + m_2)} + \frac{(m_2 p_1 - m_1 p_2)^2}{2(m_1 + m_2)^2 m_1 m_2}$$

$$= \frac{1}{2(m_1 + m_2)} \left[p_1^2 + 2p_1 p_2 + p_2^2 + \frac{m_2^2 p_1^2 - 2m_1 m_2 p_1 p_2 + m_1^2 p_2^2}{m_1 m_2} \right]$$

$$= \frac{1}{2(m_1 + m_2)} \left[p_1^2 \left(1 + \frac{m_2}{m_1}\right) + p_2^2 \left(1 + \frac{m_1}{m_2}\right) \right]$$

$$= \frac{p_1^2}{2m_1} + \frac{p_2^2}{2m_2}. \quad \text{Thus}$$

$$H = \frac{\vec{p}_1^2}{2m_1} + \frac{\vec{p}_2^2}{2m_2} + V(\vec{r}_1 - \vec{r}_2) = \frac{P^2}{2M} + \frac{p^2}{2\mu} + V(\vec{r}).$$

5. Using the same notation as in the last problem and the definition

$$\mathbf{R} = \frac{m_1 \mathbf{r}_1 + m_2 \mathbf{r}_2}{m_1 + m_2} \quad (7)$$

of the center of mass, show that the total orbital angular momentum of the two bodies

$$\mathbf{L} = \mathbf{L}_1 + \mathbf{L}_2 = \mathbf{r}_1 \times \mathbf{p}_1 + \mathbf{r}_2 \times \mathbf{p}_2 \quad (8)$$

can be written as

$$\mathbf{L} = \mathbf{R} \times \mathbf{P} + \mathbf{r} \times \mathbf{p}. \quad (9)$$

$$\mathbf{R} \times \mathbf{P} = \left(\frac{m_1 \vec{r}_1 + m_2 \vec{r}_2}{m_1 + m_2} \right) \times (\mathbf{p}_1 + \mathbf{p}_2)$$

$$\mathbf{r} \times \mathbf{p} = (\vec{r}_1 - \vec{r}_2) \times \frac{(m_2 \mathbf{p}_1 - m_1 \mathbf{p}_2)}{(m_1 + m_2)}$$

So

$$\vec{R} \times \vec{P} + \vec{r} \times \vec{p} = \mathbf{r}_1 \times \mathbf{p}_1 \frac{m_1}{M} + \mathbf{r}_2 \times \mathbf{p}_2 \frac{m_2}{M}$$

$$+ \frac{m_1}{M} \mathbf{r}_1 \times \mathbf{p}_2 + \frac{m_2}{M} \mathbf{r}_2 \times \mathbf{p}_1 + \frac{m_2}{M} \mathbf{r}_1 \times \mathbf{p}_1 + \frac{m_1}{M} \mathbf{r}_2 \times \mathbf{p}_2$$

$$- \frac{m_2}{M} \mathbf{r}_2 \times \mathbf{p}_1 - \frac{m_1}{M} \mathbf{r}_1 \times \mathbf{p}_2$$

$$= \vec{r}_1 \times \vec{p}_1 + \vec{r}_2 \times \vec{p}_2$$