

Path Integrals and the WKB Approximation

When a classical trajectory

$$\delta S = 0$$

is available, the leading term is proportional to $e^{iS/\hbar}$.

Assume a particle of mass m under the influence of a potential $V(x)$ is in a stationary state of energy E

$$H|E\rangle = E|E\rangle$$

$$H = \frac{p^2}{2m} + V(x)$$

$$L = \frac{m}{2} \dot{x}^2 - V(x).$$

The action S is

$$S = \int dt \left[\frac{m}{2} \dot{x}^2 - V(x(t)) \right].$$

Since the particle has energy E ,

$$\frac{m}{2} \dot{x}^2 + V(x) = E.$$

So

$$-V(x) = \frac{m}{2} \dot{x}^2 - E$$

and the action is

$$S = \int dt (m \dot{x}^2 - E)$$

$$= \int \dot{x} dt m \dot{x} - \int E dt$$

$$= \int dx p - \int E dt = \int p dx - Et.$$

Now p itself is related to E and to $V(x)$

since

$$\frac{p^2}{2m} + V(x) = E$$

So

$$\frac{p^2}{2m} = E - V(x)$$

$$p = \sqrt{2m(E - V(x))}$$

Thus the action S is

$$S = \int p dx - Et$$

$$= \int \sqrt{2m(E - V(x))} dx - Et,$$

which is called Hamilton's principal function.

So we expect to be able to approximate the wave function as

$$\langle x, t | E \rangle \propto e^{iS/\hbar} \\ = c e^{\frac{i}{\hbar} \left(\int \sqrt{2m(E-V(x))} dx - Et \right)}$$

where c is some factor.

The function $c(x)$ is less important than the exponential; in section (2.4) $c(x)$ is shown to be proportional to

$$c(x) \propto [E - V(x)]^{-1/4}$$

The full WKB approximation then is

$$\langle x, t | E \rangle = \frac{N}{[E - V(x)]^{1/4}} e^{\frac{i}{\hbar} \left(\int dx \sqrt{2m(E - V(x))} - Et \right)}$$

It works best in the limit $\hbar \rightarrow 0$ and when the potential $V(x)$ changes little over a wavelength λ

$$\lambda = \frac{h}{\sqrt{2m(E - V(x))}} \ll \frac{4\pi(E - V(x))}{|dV/dx|} \quad (2.4.37)$$

compared to the length scale of the problem.