

23 March 2023

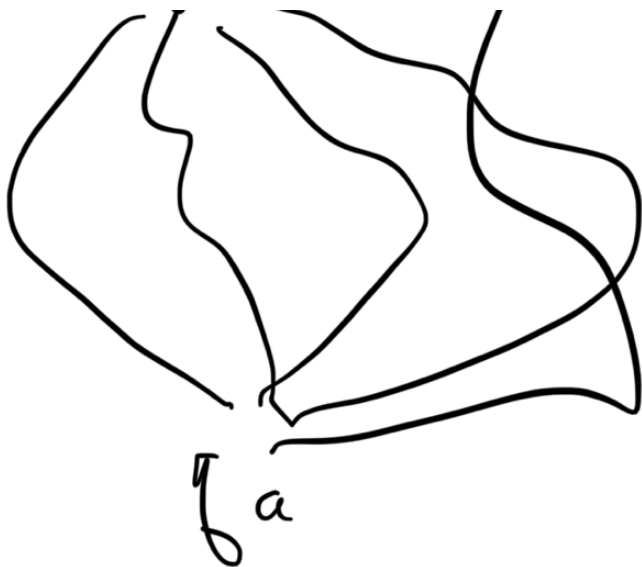


$$\beta = \frac{1}{kT}$$

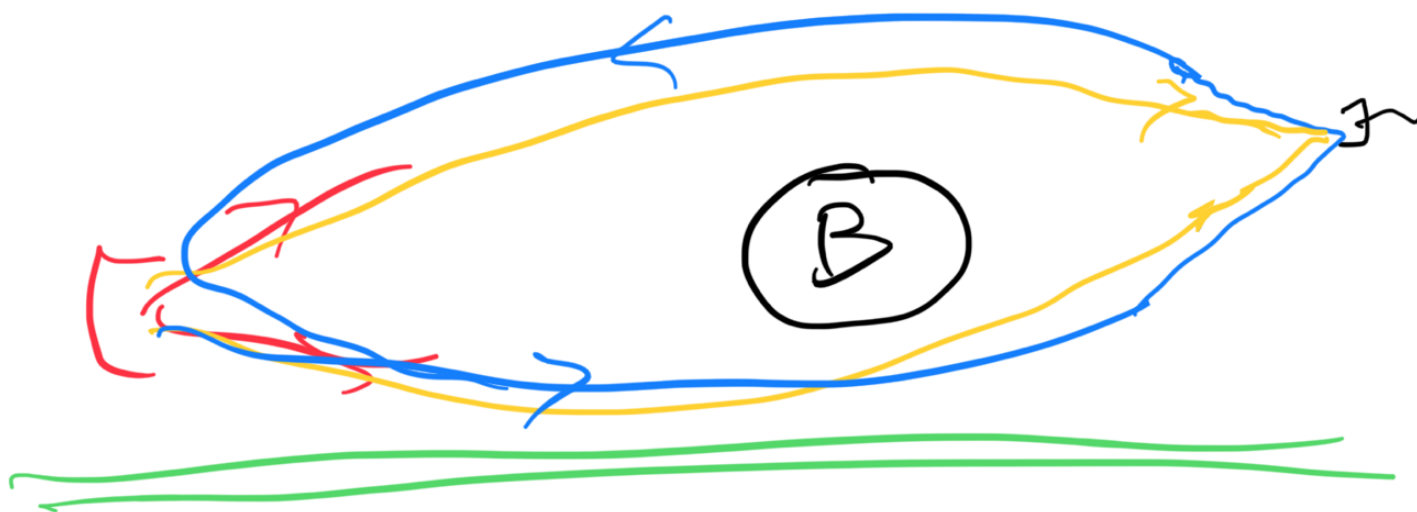
$$S = \int (T - V) dt$$

$$e^{-i\epsilon k} e^{-i\epsilon v} I e^{i\epsilon k} e^{i\epsilon v}$$

for



$$S = -mc^2 \int \sqrt{1 - v^2/c^2} dt$$



$$\beta = \frac{1}{kT}$$

$$L = T - V$$

∂^2 $\backslash / \backslash \backslash$

$$= \frac{1}{2m} - V(q)$$

$$S = \int L dt$$

$$\langle \quad \rangle = \int e^{iS[q]} Dq$$

$$\langle \{q\}_b | e^{-itH/\hbar} | \{q\}_a \rangle$$

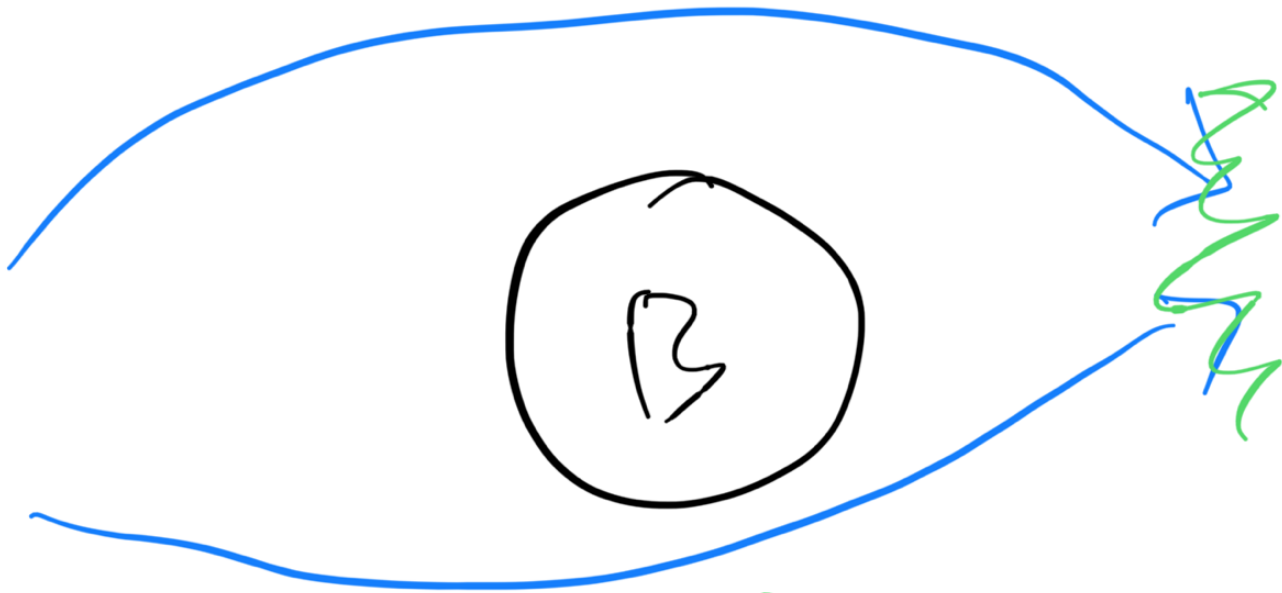
$$= \int e^{\frac{i}{\hbar} \left(\int dt \left[\frac{1}{2} m \dot{q}_j^2 - V(\{q\}) \right] \right)} Dq$$

$$V(q) = V(q_0) + V'(q_0)(q - q_0)$$

$$+ \frac{1}{2} V''(q_0) (q - q_0)^2$$

$$\left[+ \frac{1}{6} V'''(q_0) (q - q_0)^3 + \dots \right]$$

small



$$\int A \cdot dg = \int (\nabla \times A) \cdot d\mathbf{r}$$
$$= \int B \cdot d\mathbf{s} = \Phi$$

$$\text{Tr}(AB) = \text{Tr}(BA)$$

$$\text{Tr}(e^{-\beta H}) = \text{Tr}(u u^\dagger e^{-\beta H})$$

$$\text{Tr}(u^\dagger e^{-\beta H} u)$$

$$\sum_n \langle n | e^{-\beta H} | m \rangle = \sum_\alpha \langle \alpha | e^{-\beta H} | \alpha \rangle$$

$$|\alpha\rangle = U |m\rangle$$

$$Z(\beta) = \sum_n e^{-\beta E_n}$$

$$\int \boxed{g_m} d f_m$$

$$\beta = \frac{1}{kT}$$

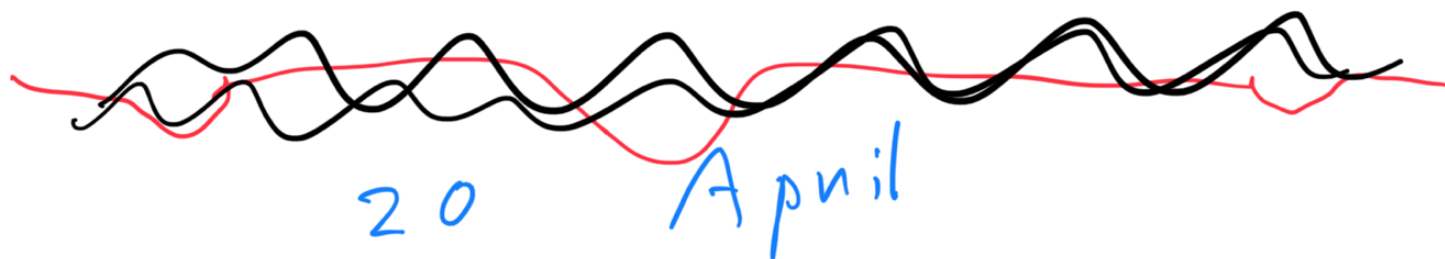
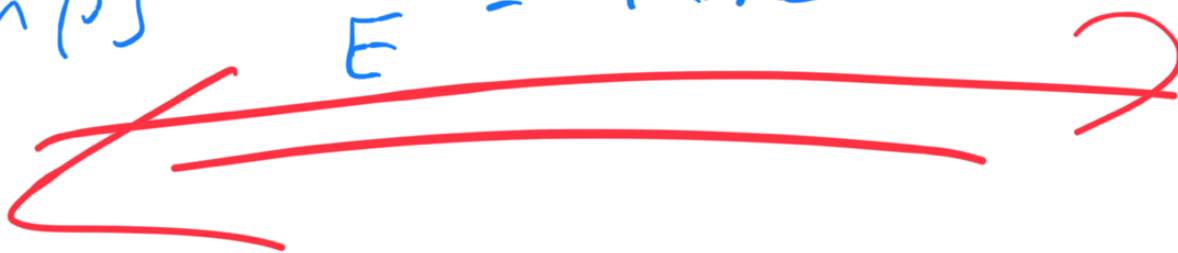
T temp.

$$T \rightarrow 0 \Leftrightarrow \beta \rightarrow \infty$$

$$\hbar \beta = m \frac{\dot{x}^2}{2} \quad \dots \quad k \rightarrow \infty \quad \uparrow$$

$$\int_0^{\infty} du \left(\frac{m v_i^2}{2} + V(v_i) \right)$$

$$[h \beta] = \frac{E T}{E} = \text{Time}$$



$$\int f(x) \delta'(x) dx$$

$$= \int -f'(x) \delta(x) dx$$

$$= -f'(0)$$

$$\begin{aligned}
 & \int f(x) \delta''(x-y) dx \\
 &= \int f''(x) \delta(x-y) dx \\
 &= f''(y).
 \end{aligned}$$

$$\{Z[j][\phi_0][\phi]\}$$

$$\frac{d}{dx} \frac{y'}{\sqrt{1+y'^2}} = \frac{y'}{\sqrt{1+y'^2}} - \frac{1}{2} \frac{y'' \cdot 2y'y'}{(1+y'^2)^{3/2}} = 0$$

$$y'' = 0 = \frac{y''}{\sqrt{1+y'^2}} \left(1 - \frac{y'y''}{1+y'^2} \right)$$

$$\binom{a}{b} = \frac{a!}{b!(a-b)!}$$

$$\langle q' | f \rangle = f(q')$$

$$\langle q' | \psi \rangle = \Psi(q')$$

$$\begin{aligned} \langle q' | \underline{pq} | f \rangle &= \langle q' | \frac{\hbar}{i} + q' p | f \rangle \\ &= \langle q' | \underline{\frac{\hbar}{i} + qp} | f \rangle \end{aligned}$$

$$[q, p] = i\hbar \quad qp - pq = i\hbar$$

$$[\phi(\vec{x}, t), \phi(\vec{y}, t)] = 0$$

$$\phi(x, t) |\psi'\rangle = \underset{\text{real}}{\phi'(x)} |\psi'\rangle$$

$$\underset{\text{hermitian}}{\langle \phi' | \phi'(x)} = \langle \phi' | \phi^\dagger(x, t) = \langle \phi' | \phi(x, t)$$

$$\hbar = 1 = c$$

$$-\sqrt{-\nabla^2 + m^2} \phi'(x) \langle \phi' | 0 \rangle \quad \langle \phi' | 0 \rangle$$

$$= -\int \delta(x-x') \sqrt{-\nabla'^2 + m^2} \phi(x') d^3x'$$

$$= -\sqrt{-\nabla^2 + m^2} \phi'(x) \langle \phi' | 0 \rangle$$

$$\sqrt{-\nabla^2 + m^2} \int e^{+i\vec{p}\cdot\vec{x}} \tilde{\phi}(p) \frac{d^3p}{(2\pi)^{3/2}}$$

$$= \int \sqrt{p^2 + m^2} e^{ip\cdot x} \tilde{\phi}(p) \frac{d^3p}{(2\pi)^{3/2}}$$

$$\left(\int d^3x \int \frac{d^3p d^3p'}{(2\pi)^3} \tilde{\phi}(p') e^{ip'\cdot x} \tilde{\phi}(p) e^{ip\cdot x} \right. \\ \left. \times \sqrt{p^2 + m^2} \right)$$

$$\int d^3p d^3p' \sqrt{p^2 + m^2} \delta(\vec{p} + \vec{p}')$$

$$= \int d^3p \tilde{\phi}(p') \tilde{\phi}(p)$$

$$= \int d^3p \tilde{\phi}(-p) \tilde{\phi}(p) \sqrt{p^2 + m^2}$$

$$= \int d^3p \tilde{\phi}^*(p) \tilde{\phi}(p) \sqrt{p^2 + m^2}$$

19.54

$$\langle \phi' | 0 \rangle = N e^{-\frac{1}{2} \int |\tilde{\phi}(\vec{p})|^2 \sqrt{\vec{p}^2 + m^2} d^3p}$$

$\phi'(x)$

27 April

$$= - \frac{\partial f(x)}{\partial y}$$

$$f(x) \delta'(x-y) \rightarrow - \delta(x-y) \frac{\partial f(x)}{\partial x}$$

$$f(x) \delta''(x-y) \rightarrow \delta(x-y) \frac{\partial^2 f(x)}{\partial x^2}$$

$$= + \frac{\partial^2 f(x)}{\partial x^2}$$

$$0 = H |0\rangle = \frac{1}{2} \int [\pi] [\sqrt{-\nabla^2 + m^2} \phi + i\pi] d^3x |0\rangle$$

$$\langle \phi | [\sqrt{-\nabla^2 + m^2} \phi + i\pi] |0\rangle = 0$$

$$\langle \phi | \pi |0\rangle = \frac{1}{i} \frac{\delta \langle \phi |0\rangle}{\delta \phi}$$

$$* \quad \langle g | p | \psi \rangle = \frac{1}{i} \frac{d}{dg} \langle g | \psi \rangle$$

$$\underline{\phi(\vec{x}, 0)} | \underline{\phi'} \rangle = \underline{\phi'(\vec{x})} | \underline{\phi'} \rangle$$

any of many commuting fields
 $\dots \phi(\vec{x}_k, 0) \dots \phi(\vec{x}_n, 0) \dots$

$$\partial^a e^{i p' x} = \frac{\partial}{\partial x_a} e^{i p' x}$$

$$= i \phi'^a$$

$$\phi(\vec{x}, 0)$$

Field

$$\phi(\vec{x}, 0) | \phi'(\vec{x}, t) \rangle$$

$$= \phi'(\vec{x}, t) | \phi'(\vec{x}, t) \rangle$$

$$(x+a)^2 = x^2 + 2xa + a^2$$

$$(x+a)^2 - a^2 = x^2 + 2xa$$

$$(x+a)^2 b^2 = x^2 b^2 + 2xab^2 + a^2 b^2$$

$$(x+a)^2 b^2 - a^2 b^2 = x^2 b^2 + 2xab^2$$

$$\left(\frac{1}{i} \right)^2 \frac{\delta}{\delta i(x)} \frac{\delta}{\delta i(x')} Z_0[J] \Big|_{\hat{i}=0}$$

$$= \langle 0 | T (\phi(x) \phi(x')) | 0 \rangle$$

4 May 2023

local $\Leftrightarrow$ depends upon
 spacetime $x = (ct, \vec{x}) \Leftrightarrow$ gauge

$$\psi' = \begin{pmatrix} \psi_1 \\ \psi_2 \\ \vdots \\ \psi_n \end{pmatrix} = U \begin{pmatrix} \psi(x) \\ \psi(x) \end{pmatrix}$$

$n \times n$

$\psi' = U \psi$

$$(\psi_{(x)}^\dagger \psi_{(x)})' = \psi^\dagger U^\dagger U \psi = \psi^\dagger \psi$$

$$U^\dagger U = I$$

$$[(D_i \psi)^\dagger \gamma^i \psi]' = (D_i \psi)^\dagger \gamma^i \psi$$

$$(\psi^\dagger D_i \gamma^i \psi)'$$

$$= \psi^\dagger u^\dagger (D_i \gamma^i \psi)'$$

$$= \psi^\dagger u^\dagger u D_i \gamma^i \psi$$

$$= \psi^\dagger D_i \gamma^i \psi$$

$$\{\gamma^a, \gamma^b\} = \gamma^a \gamma^b + \gamma^b \gamma^a = 2\eta^{ab}$$

$$= 2 \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$U^\dagger \rightleftharpoons \tilde{U}^{-1}$$

$$U^\dagger = \tilde{U}^{-1}$$

$$i \gamma^0(x)$$

electrodynamics $U(x) = e^{\nu}$

$$[t_a, t_b] = t_a t_b - t_b t_a$$

$$= i f_{ab}^c t_c$$

↙ structure constants

$$A'_i(x) = A_i(x) + \partial_i \lambda(x)$$

QED

$$D'_i = U D_i U^\dagger = U D_i U^\dagger$$

$$F_{ik} = (\partial_i + A_i)(\partial_k + A_k) - (\partial_k + A_k)(\partial_i + A_i)$$

$$= \partial_i A_k - \partial_k A_i + A_i A_k - A_k A_i$$

$$= \partial_i A_k - \partial_k A_i + [A_i, A_k]$$

$\uparrow \quad \nearrow$
 $m \times m$

$$\vec{P} = \frac{\hbar}{i} \nabla$$

herm. antiherm.

$$\psi = \begin{pmatrix} l \\ r \end{pmatrix}$$

$z=5$
 e

$z=5$
 e

$$\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$L_e = \begin{pmatrix} \nu \\ e \end{pmatrix} \quad \begin{array}{l} \text{neutrino} \\ \text{electron} \end{array}$$

$$SU_2(2)$$

$$D_i L_e = (\partial_i + A_i) L_e$$

$$= \left(\partial_i - ig \frac{\vec{\sigma}}{2} \cdot \vec{A}_i \right) L_e$$

$$U_Y(1)$$

$$D_i L_e = (\partial_i - ig' Y B_i) L_e$$

$$D_i R_e = (\partial_i - ig' Y B_i) R_e$$

$$R_e = \begin{pmatrix} \nu_e \\ e \end{pmatrix}$$

$$\mathcal{H} = \dots \lambda^2 (\phi^\dagger \phi - \mu^2)^2$$

$$\langle 0 | \phi^\dagger \phi | 0 \rangle = \mu^2 \quad ; \phi$$

$$\langle 0 | \phi(x) | 0 \rangle = \mu e$$

$$\langle 0 | \phi(x) | 0 \rangle = \mu$$

$$S = -\frac{1}{2} \mu^2 \phi^2$$

$$\text{mass} = \mu$$

$$H = \begin{pmatrix} h^+ \\ h^0 \end{pmatrix} =$$

$$\langle 0 | H | 0 \rangle = \begin{pmatrix} 0 \\ v \end{pmatrix}$$

$$(D; \psi)^{\dagger} (D; \psi)$$

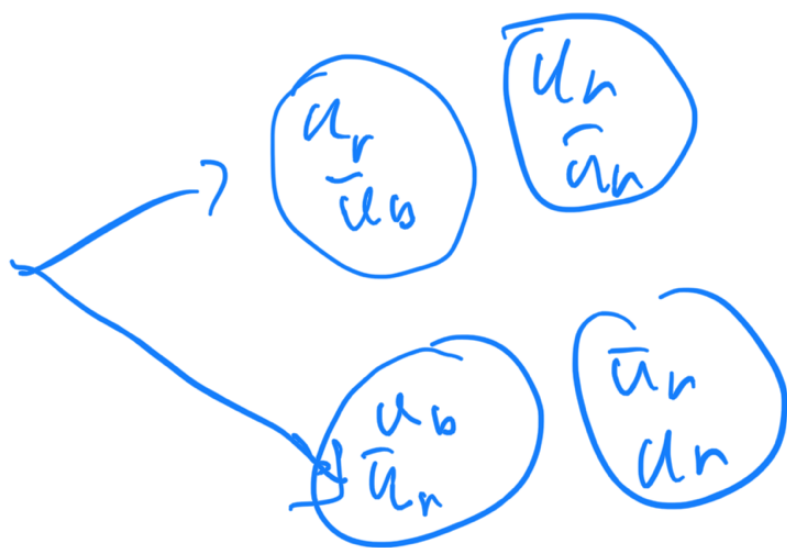
$$(D; H)^{\dagger} (D; H)$$

$$\vec{u} = \begin{pmatrix} \vec{l}_u \\ \vec{r}_u \end{pmatrix} = \begin{pmatrix} \psi_u \\ \psi_u \\ \psi_u \end{pmatrix} = \mathbb{I}_u$$

$$SU_c(3)$$

$$D_i \Psi_u = \partial_i \bar{\Psi}_u + A_i \bar{\Psi}_u$$

$$A_i = -ig \lambda^\alpha A_i^\alpha$$



$$\underline{\underline{A_i(x)}} \rightarrow \underline{\underline{V(x)}}$$



