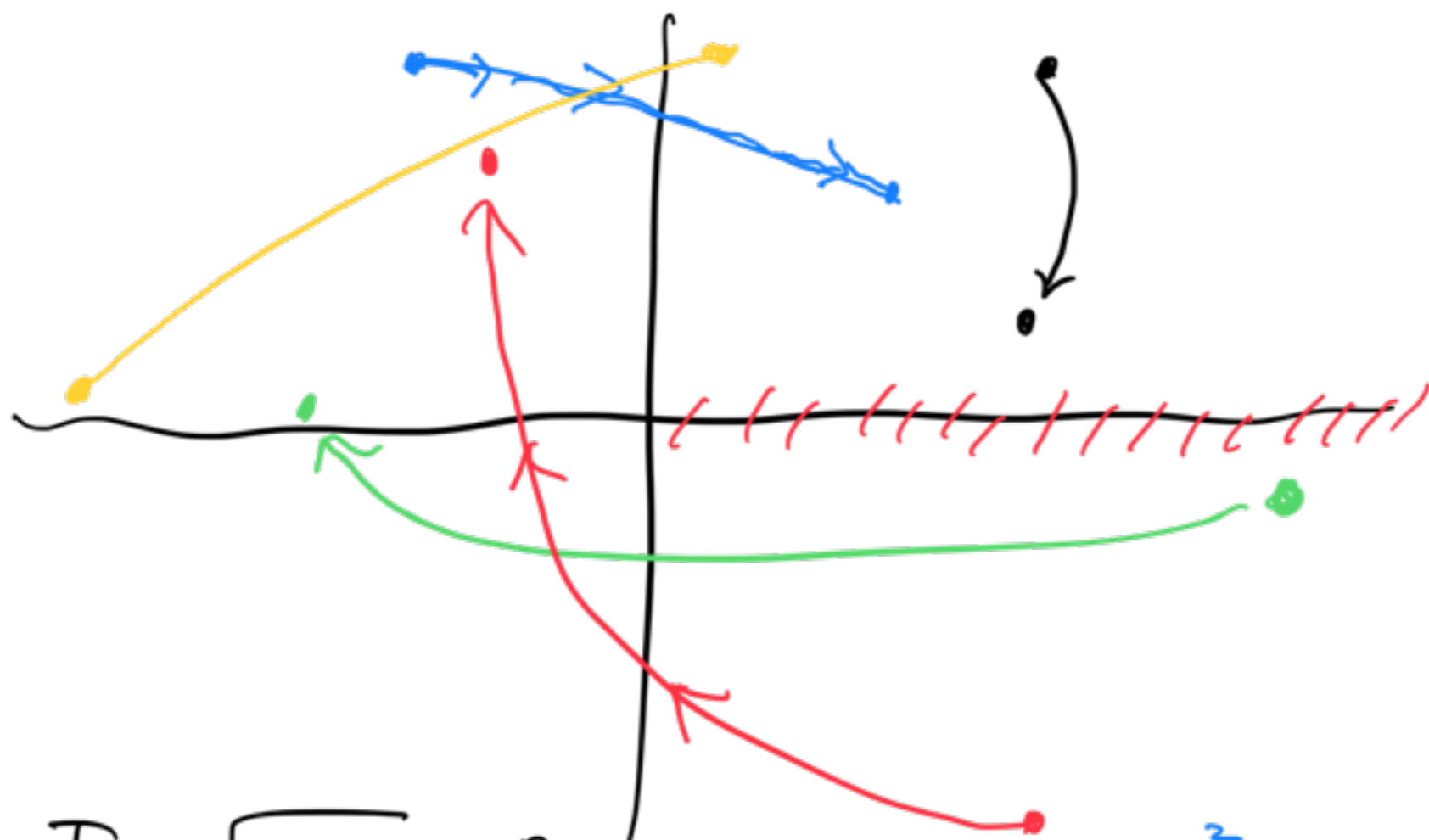
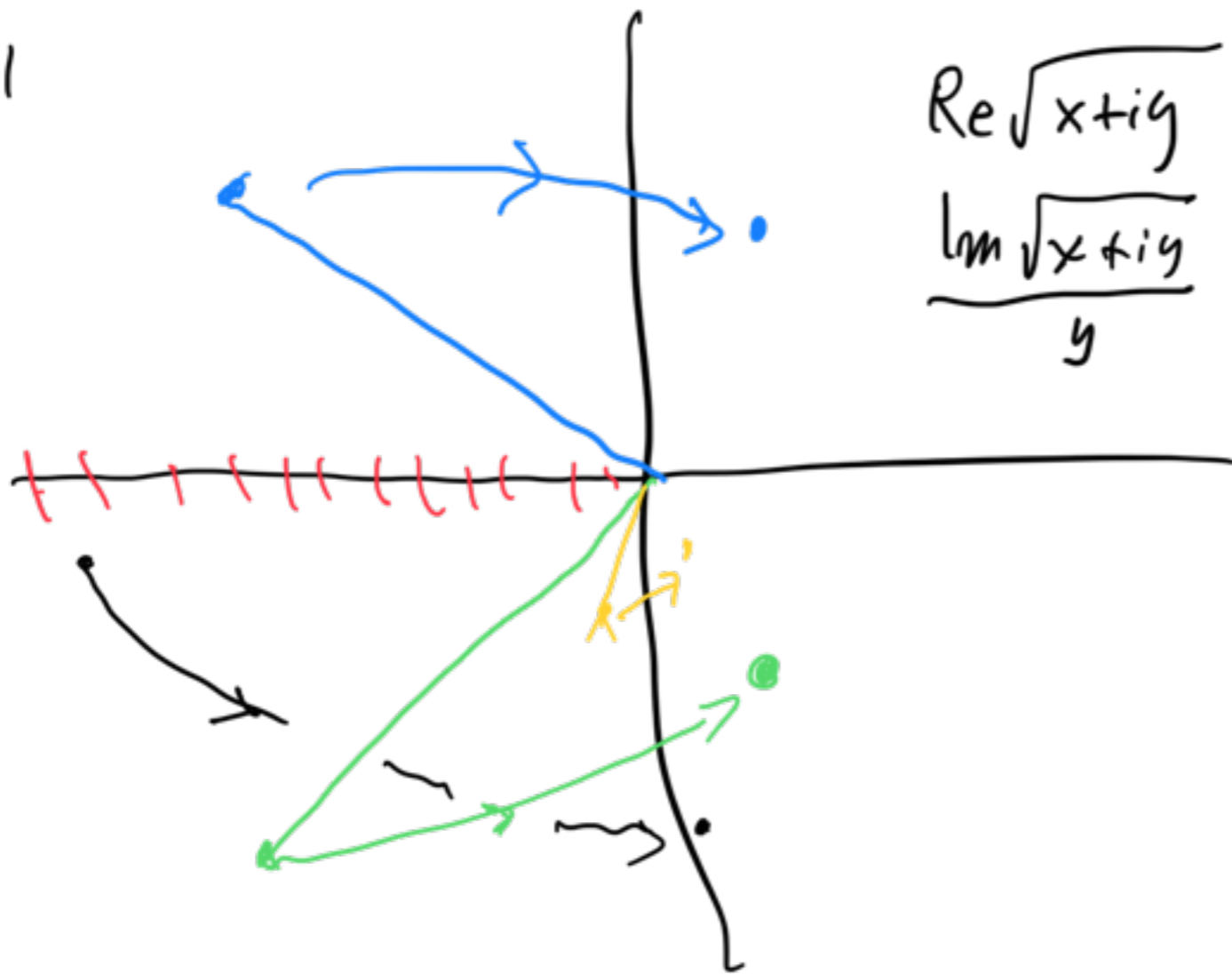


6.31

$$\operatorname{Re} \sqrt{x+iy} > 0$$

$$\frac{\operatorname{Im} \sqrt{x+iy}}{y} > 0$$



$$\operatorname{Im} \sqrt{x+iy} > 0$$

$$x+iy = \left( \frac{\alpha}{|\alpha|} + i\beta \right)^2 = \frac{\alpha^2 - \beta^2}{|\alpha|^2 + \beta^2} + i \frac{2\alpha\beta}{|\alpha|^2 + \beta^2}$$

$$\frac{\operatorname{Re} \sqrt{x+iy}}{y} > 0$$

$$\sqrt{x+iy} = \frac{y}{|y|} \alpha + i\beta$$



$$\sqrt{x+iy} = \alpha + i\beta \frac{y}{|y|}$$

$$(x+iy) = \alpha^2 - \beta^2 + 2i\alpha\beta \frac{y}{|y|}$$

$$x = \alpha^2 - \beta^2$$

$$|y| = 2\alpha\beta$$

$$\alpha^2 = x + \beta^2 = x + \frac{y^2}{4\alpha^2}$$

$$\alpha^4 - x\alpha^2 - \frac{y^2}{4} = 0$$

$$\alpha^2 = \frac{x + \sqrt{x^2 + y^2}}{2}$$

$$0 < \alpha^2 \Rightarrow \alpha = \sqrt{\frac{x + \sqrt{x^2 + y^2}}{2}}$$

$$\beta^2 = \alpha^2 - x = \frac{x + \sqrt{x^2 + y^2}}{2} - x = \frac{\sqrt{x^2 + y^2} - x}{2}$$

$$\beta = \sqrt{\frac{\sqrt{x^2 + y^2} - x}{2}}$$

$$i = e^{i\frac{\pi}{2}}$$

$$i = e^{i\frac{\pi}{2}} = e^{2\pi ni}$$

$i = e^{i(\frac{\pi}{2} + 2\pi n)}$   $n$  an integer

$i = e^{i(\frac{\pi}{2} + 2\pi n)}$

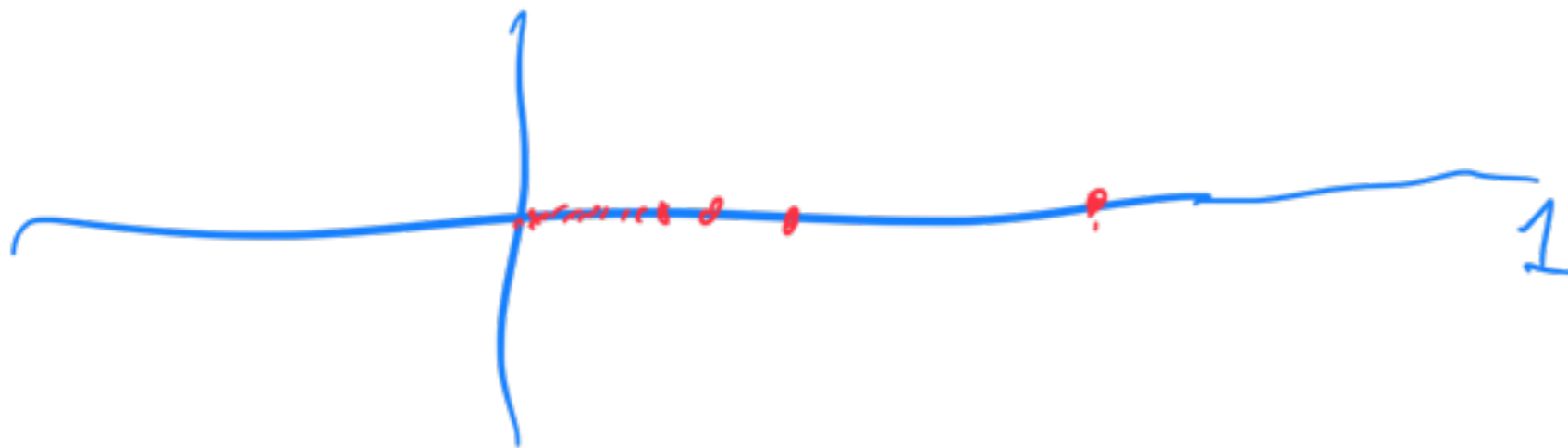
$i^i \neq e^{-\frac{\pi}{2} - 2\pi n}$

$i^i = e^{i \cdot i(\frac{\pi}{2} + 2\pi n)} = e^{-\frac{\pi}{2} - 2\pi n}$

$i^i = e^{i(\frac{\pi}{2} + 2\pi n) e^{i(\frac{\pi}{2} + 2\pi n)}}$

$i^i = e^{-\frac{\pi}{2} - 2\pi n}$

$i^i = e^{i^2(\frac{\pi}{2} + 2\pi n)} = e^{-\frac{\pi}{2} - 2\pi n}$



7.8



$$V = IR + \frac{Q}{C}$$

$$\dot{V} = I\dot{R} + \frac{I}{C}$$

$$\dot{I} + \frac{I}{RC} = \frac{\dot{V}}{R}$$

$$r = \frac{1}{RC} \quad s = \frac{\dot{V}(t)}{R}$$

$$\alpha(t) = \exp\left(\int \frac{dt'}{RC}\right) = e^{\frac{t}{RC}}$$

$$y(t) = e^{-\frac{t}{RC}} \left[ I(0) + \int_0^t e^{\frac{t'+t}{RC}} \frac{\dot{V}(t')}{R} dt' \right]$$

$$= e^{-\frac{t}{RC}} \frac{1}{R} \int_0^t e^{\frac{t'+t}{RC}} \dot{V}(t') dt'$$

$$\dots \cos(\omega t')$$

$$V(t') = V_0 \cos(\omega t')$$

$$= e^{-\frac{t}{RC}} \frac{V_0}{R} \int_0^t e^{t'/RC} \cos(\omega t') dt'$$

$$= \frac{\omega V_0}{2R} e^{-\frac{t}{RC}} \int_0^t \left( e^{t'(\frac{1}{RC} + i\omega)} + e^{t'(\frac{1}{RC} - i\omega)} \right) dt'$$

$$= \frac{\omega V_0}{2R} e^{-\frac{t}{RC}} \left[ \frac{e^{t(\frac{1}{RC} + i\omega)}}{\frac{1}{RC} + i\omega} + \frac{e^{t(\frac{1}{RC} - i\omega)}}{\frac{1}{RC} - i\omega} \right]$$

$$= \frac{\omega V_0 C}{2(1 + \omega^2 R^2 C^2)} \left[ -2e^{-\frac{t}{RC}} + 2\cos \omega t + 2\omega R C \sin \omega t \right]$$

$$= \frac{\omega V_0 C}{1 + \omega^2 R^2 C^2} \left( \cos \omega t + \omega R C \sin \omega t - e^{-\frac{t}{RC}} \right)$$

$$\dot{y} = -e^{-t} \quad y(0) = 1$$

$$y(t) = \log(e - t)$$

$$\dot{y} = e^{-y} \quad y = \log(-t + e)$$

Ex. 7.45 Hermite polynomials

$$y'' - x^2 y + \lambda y = 0 \quad -\infty < x < \infty$$

↑ diverges as  $x \rightarrow \infty$  is essential singularity

For large  $x$   $y'' \approx x^2 y$

$$\text{So } y(x) \sim e^{-x^2/2} \quad y' = -x e^{-x^2/2}$$

$$y'' = x^2 e^{-x^2/2} - e^{-x^2/2} \text{ so } y'' \approx x^2 y$$

for large  $x$ . So we set

$y(x) = e^{-x^2/2} h(x)$  and get

$$y' = -x y + h' e^{-x^2/2} = e^{-x^2/2} (h' - x h)$$

$$y'' = e^{-x^2/2} (-x h' + x^2 h + h'' - h - x h')$$

$$= e^{-x^2/2} (h'' - 2xh' + (\lambda-1)h)$$

$$\text{So } y'' + (\lambda - x^2)y = 0 \Rightarrow$$

$$e^{-x^2/2} [h'' - 2xh' + (\lambda-1)h + (\lambda - x^2)h] = 0$$

$$h'' - 2xh' + (\lambda-1)h = 0 \quad \text{So}$$

$$\text{we let } h(x) = x^r \sum_{n=0}^{\infty} a_n x^n = \sum_{n=0}^{\infty} a_n x^{n+r}$$

$$\text{So } h' = \sum a_n (n+r) x^{n+r-1}$$

$$h'' = \sum (n+r)(n+r-1) a_n x^{n+r-2}$$

$$\sum (n+r)(n+r-1) a_n x^{n+r-2} - \sum 2(n+r) x^{n+r-1} a_n + \sum (\lambda-1) a_n x^{n+r} = 0$$

$$r(r-1) = 0 \quad \text{so } r=0 \text{ \& } r=1$$

For  $r=0$

$$0 = \sum n(n-1) a_n x^{n-2} - \sum 2n x^{n-1} a_n + \sum (\lambda-1) a_n x^n$$

$$0 = \sum \{(n+2)(n+1) a_{n+2} + [-2n + \lambda - 1] a_n\} x^n = 0$$

$$(n+2)(n+1) a_{n+2} = (2n+1-\lambda) a_n$$

$$a_{n+2} = \frac{2n+1-\lambda}{(n+2)(n+1)} a_n$$

Unless  $\lambda = 2n+1$ , the

series is like  $a_{n+2} \sim \frac{2}{n+2} a_n$

$$\text{or } a_n \sim \frac{2}{n} a_{n-2}$$

$$\text{or } a_{2m} \sim \frac{2^m a_0}{(2m)!!}$$

$$\text{So } h(x) \sim \sum a_{2m} x^{2m} \sim \sum \frac{2^m x^{2m}}{(2m)!!} a_0$$

$$\sim \sum \frac{(x^2)^m}{m!} a_0 \sim e^{x^2}$$

$$\text{But then } y(x) \sim e^{-x^2/2} e^{x^2} \\ \sim e^{x^2/2}$$

is not square integrable.



So  $\lambda = 2n+1$  and  
then  $h(x)$  is a polynomial,  
a Hermite polynomial.

---

For  $r=1$

$$a_{n+2} = \frac{2n+1-\lambda}{(n+2)(n+1)} a_n$$

and again

$\lambda = 2n+1$   
is an eigenvalue.

---

Here  $y'' + (\lambda - x^2)y = 0 \quad \square$

$$y'' + Q(x)y = 0$$

$$P(x) = 0$$

$$Q(x) = \lambda - x^2$$

$$z = \frac{1}{x}$$

$$\frac{Q(1/2)}{-4} = \frac{\lambda - 1/2^2}{-4} = \frac{\lambda}{-4} - \frac{1}{2^6}$$

diverges as  $z \rightarrow 0$ .

And  $\frac{Q(z)}{z^2} = \frac{\lambda}{z^2} - \frac{1}{z^4}$  also diverges

So  $x = \pm\infty$  are essential singularities of the Hermite equation.

---

H atom

$$(r^2 R')' + (\alpha r^2 + \beta r + \gamma) R = 0$$

$$r^2 R'' + 2r R' + (\alpha r^2 + \beta r + \gamma) R = 0$$

$$R'' + \frac{2}{r} R' + \left(\alpha + \frac{\beta}{r} + \frac{\gamma}{r^2}\right) R = 0$$

$$P(r) = \frac{2}{r} \quad Q(r) = \alpha + \frac{\beta}{r} + \frac{\gamma}{r^2}$$

$r=0$  is a singular point

But  $rP(r) = 2$  and  $r^2Q = 2r^2 + \beta r + \gamma$

are both finite at  $r=0$ .

So  $r=0$  is a regular singular point of the non-rel. H-atom eq.

is a regular

But  $r = \infty$  is an essentially singular point.

We first set  $R = r^l S$

$$R' = l r^{l-1} S + r^l S'$$

$$r^2 R' = l r^{l+1} S + r^{l+2} S'$$

$$(r^2 R')' = l(l+1) r^l S + (l+2) r^{l+1} S' + r^{l+2} S''$$

Need to cancel  $l(l+1) r^l$  with

$$(\alpha r^2 + \beta r + \gamma) r^l S \quad \text{So } \gamma = -l(l+1).$$

Next  $S = e^{-\delta r}$  and find

$$r^2 R'' \sim r^2 \delta^2 + \alpha r^2 = 0$$

So we need  $\alpha \leq 0$  and  $\delta = \sqrt{-\alpha}$

$$\frac{d}{dx} \rightarrow \frac{d}{d(-x)} = - \frac{d}{dx}$$

$$\rightarrow \hbar^2 \quad d^2$$

$$2m \frac{d^2}{dx^2}$$

$$\hbar^2 \frac{d^2}{dx^2}$$

$$-\frac{\hbar^2}{2m} \nabla^2$$

$$V(x) = V(-x)$$

$$V(\vec{r}) = V(-\vec{r}).$$

$$L(x)g(x) = 0$$

$$L(-x)g(-x)$$

$$L(-x)g(-x) = 0$$

$$L(-x) = \pm L(x)$$

$$\pm L(x)g(-x) = 0$$

$$g(x), g(-x)$$

$$\begin{array}{ll}
 y(x) + y(-x) & \text{even} \\
 y(x) - y(-x) & \text{odd}
 \end{array}$$


---

Wronski

$$Y_{ij} = \begin{pmatrix} y_1 & y_2 \\ y_1' & y_2' \end{pmatrix}$$

$$W = |Y| = y_1 y_2' - y_2 y_1' \neq 0$$

$$W' = \underbrace{y_1' y_2'} + y_1 y_2'' \underbrace{- y_2' y_1'} - y_2 y_1''$$

$$= y_1 y_2'' - y_2 y_1''$$

$$y_1'' + P y_1' + Q y_1 = 0$$

$$y_2'' + P y_2' + Q y_2 = 0$$

$$\therefore = y_1 (-P y_2' - Q y_2) - y_2 (-P y_1' - Q y_1)$$

$$W = y_1 y_2'$$

$$= -P y_1 y_2' + P y_2 y_1'$$

$$= -P W$$

$$\frac{W'}{W} = -P = (\log W)'$$

$$\log W(x) = \int^x P(x') dx' - \int^x P(x') dx'$$

$$W(x) = W(x_0) e^{x_0}$$

$$= y_1 y_2' - y_2 y_1'$$

$$= y_1^2 \frac{d}{dx} \frac{y_2}{y_1} = y_2 y_1' - y_1 y_2 \frac{y_1'}{y_1^2}$$

$$= y_1 y_2' - y_2 y_1'$$

$$W = y_1^2 \left( \frac{y_2}{y_1} \right)' = e^{-\int P(x') dx'}$$

$$\left( \frac{y_2}{y_1} \right)' = \frac{1}{y_1^2} e^{-\int P(x') dx'}$$

$$\frac{y_2}{y_1} = \int \frac{e^{-\int p(x') dx'}}{y_1^2(x')} dx'$$

$$Y = \begin{pmatrix} y_1^{(l_1)}(x_1) & y_2^{(l_1)}(x_1) \\ y_2^{(l_2)}(x_1) & y_2^{(l_2)}(x_2) \end{pmatrix}$$

$$Y C = B$$

$$Y \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = B = \begin{pmatrix} b_1 \\ b_2 \end{pmatrix}$$

$$c_1 y_1(x_1) + c_2 y_2(x_1) = b_1$$

$$c_1 y_1(x_2) + c_2 y_2(x_2) = b_2$$

$$C = Y^{-1} B$$

$$Y = \begin{pmatrix} y_1(x_1) & y_2(x_1) \\ y_1(x_2) & y_2(x_2) \end{pmatrix}$$

$$Y_{jk} = g_k(x_j)$$

---

$$\begin{cases} \alpha \sinh a + \beta \cosh a = c \\ \alpha \sinh(-a) + \beta \cosh(-a) = d \end{cases}$$

$x = \pm a$

$$-\alpha \sinh a + \beta \cosh a = d$$

$$\beta \cosh a = d + \alpha \sinh a$$

$$2\alpha \sinh a + d = c$$

$$\alpha = \frac{c-d}{2\sinh a}$$

$$\dots (c-d)/2 \dots$$



$$\beta = \frac{a + i \dots}{\cos ka} = \frac{c + d}{2 \cos ka}$$

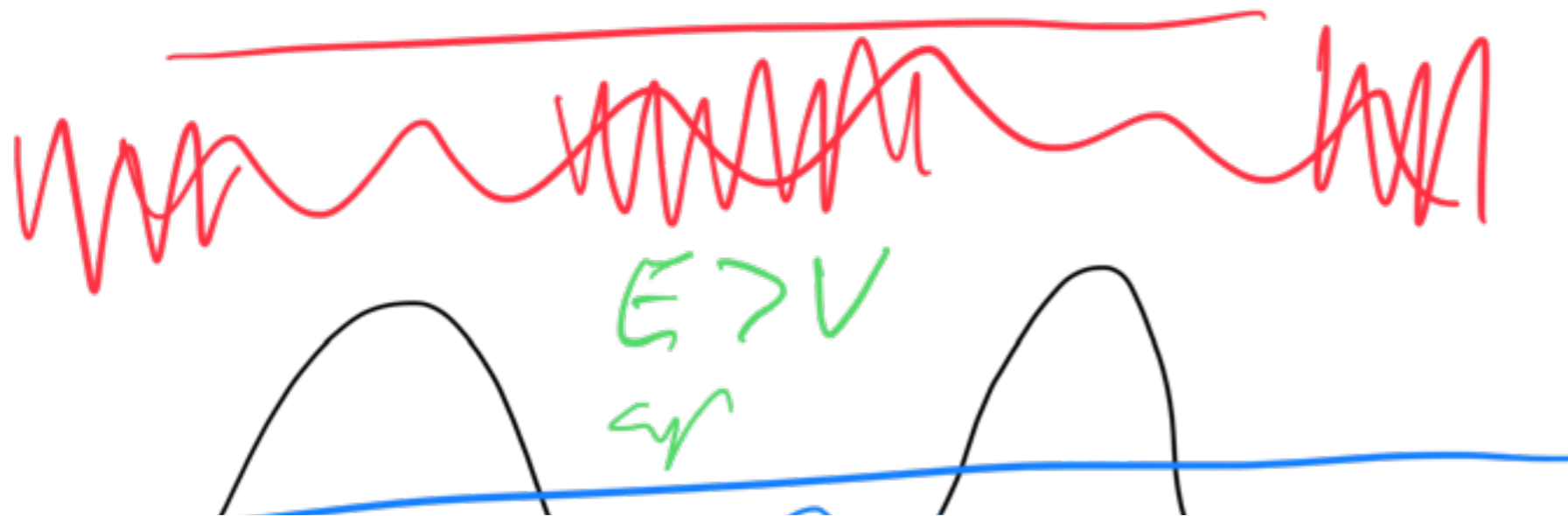
$$-y'' = k^2 y$$

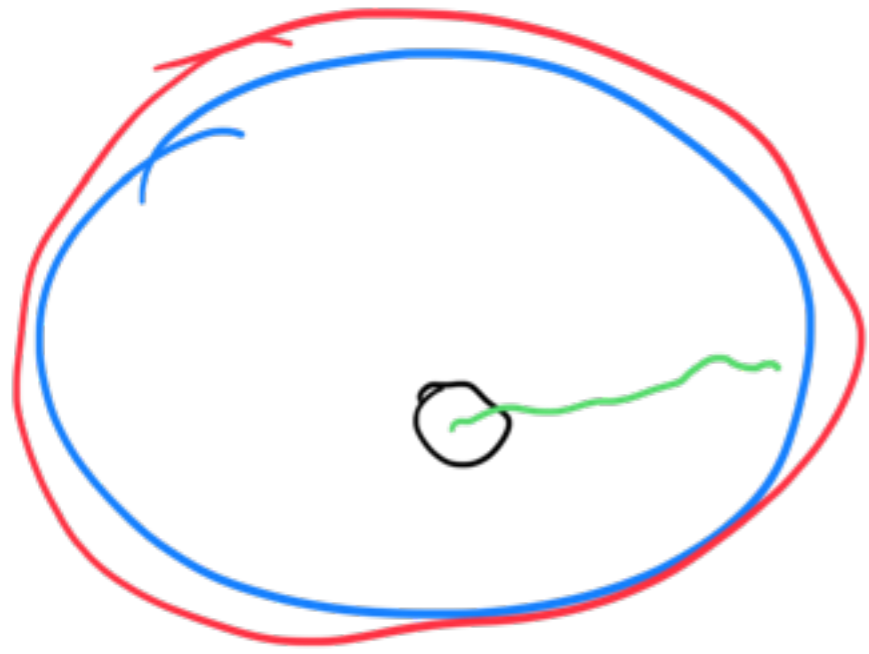
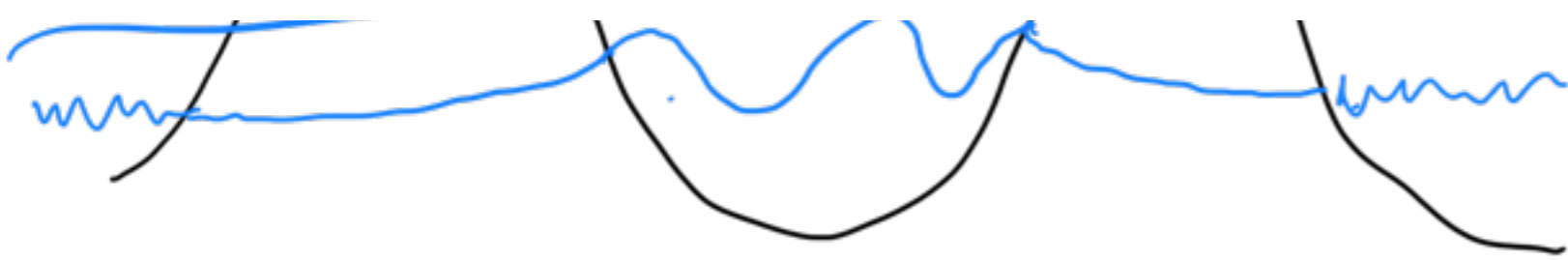
$\sin ka$        $\cos ka$

$$y(x) = \frac{c-d}{2 \sin ka} \sin hx + \frac{c+d}{2 \cos ka} \cos hx$$

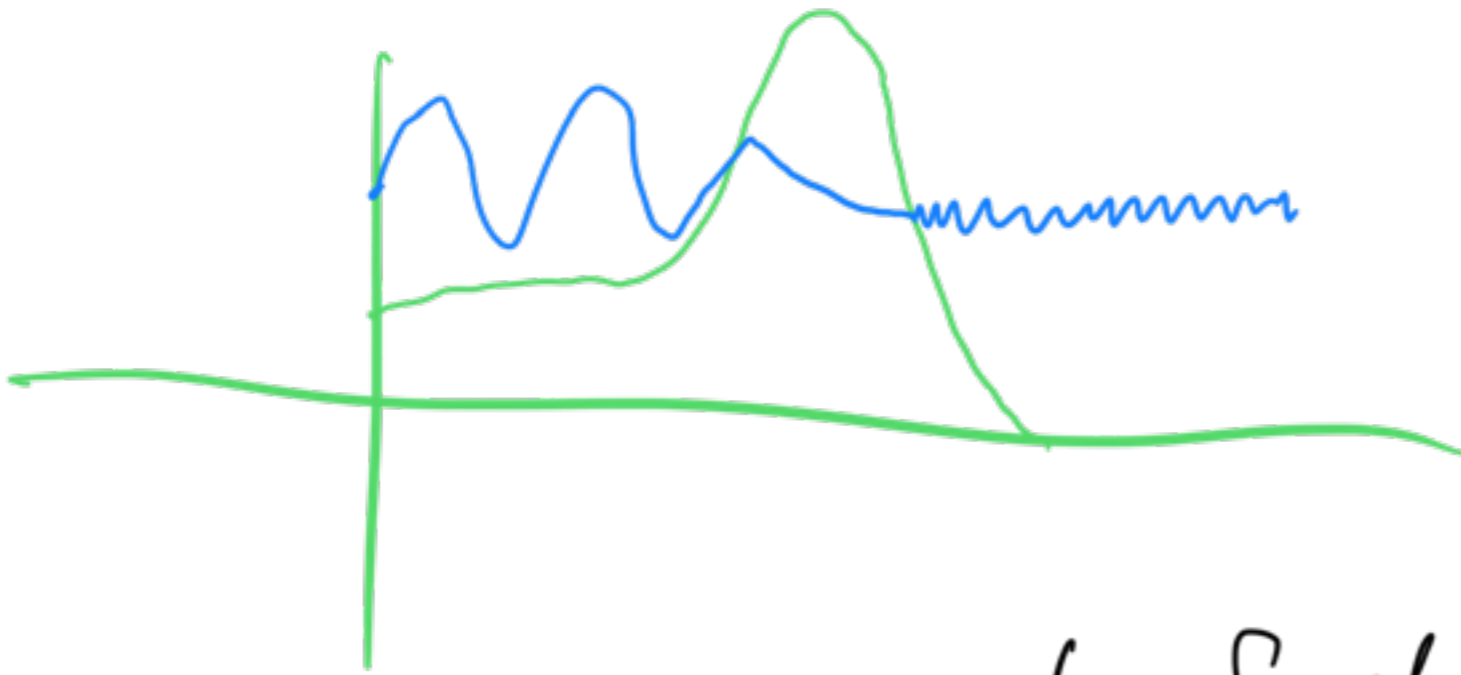
$$W_0'' = \frac{1}{2} \sqrt{\frac{2m}{\hbar^2} V'} \sqrt{-2m(E-V)/\hbar^2}$$

$$\sqrt{\frac{m V'}{\hbar^2}} + W_1'^2 = -2m \frac{(E-V)}{\hbar^2}$$





tunneling



Scale factor

$$H = \frac{\dot{a}}{a}$$

~~$$\ddot{z} + \dot{z}^2 + \left(\frac{c^2 m}{\hbar}\right)^2 = 0$$~~

$$z_0 = \frac{c^2 m}{\hbar}$$

$h$

$$\begin{aligned}\phi = e^{\pm i} &= \int_{\pm} \frac{ic^2 m}{h} - \frac{3c}{2a} dt \\ &= \left( \cos \int_{\pm} \frac{c^2 m}{h} dt' \right)^{-\frac{3}{2} \frac{t}{a}}\end{aligned}$$

---

$$L_0 = h_2 \frac{d^2}{dx^2} + h_1 \frac{d}{dx} + h_0$$

$$L_1 = -\frac{1}{h_2} L_0$$

$$= -\frac{d^2}{dx^2} - \frac{h_1}{h_2} \frac{d}{dx} - \frac{h_0}{h_2}$$

$$L = p L_1 = -p \left[ \frac{d^2}{dx^2} - \frac{h_1}{h_2} \frac{d}{dx} - \frac{h_0}{h_2} \right]$$

$$\Rightarrow \frac{e^{\int h_1/h_2}}{h_2} L_0 = p L_0$$

$$\underline{L_0 u = \lambda u}$$

$$\underline{L u = p L_0 u = p \lambda u}$$

$$\exp\left[-\int^x \frac{p'}{p} dx'\right] = \exp\left[-\int (\log p)' dx\right]$$

$$= \exp\left[-\log p(x) - c\right]$$

$$= \frac{d}{p(x)}$$

---

$$p = \frac{\hbar}{i} \frac{d}{dx} \quad \text{hermitian}$$

$$\vec{L} = \vec{x} \wedge \vec{p} = \vec{x} \times \vec{p}$$

---

$$\int u'^2 + p u^2 dx$$

$$\int u' \delta u' + p u \delta u$$

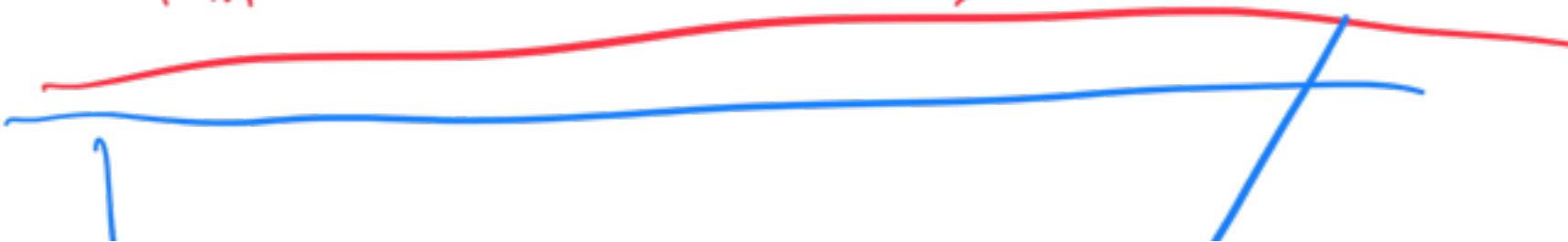
$$\Rightarrow \int -u'' \delta u + (p u) \delta u$$

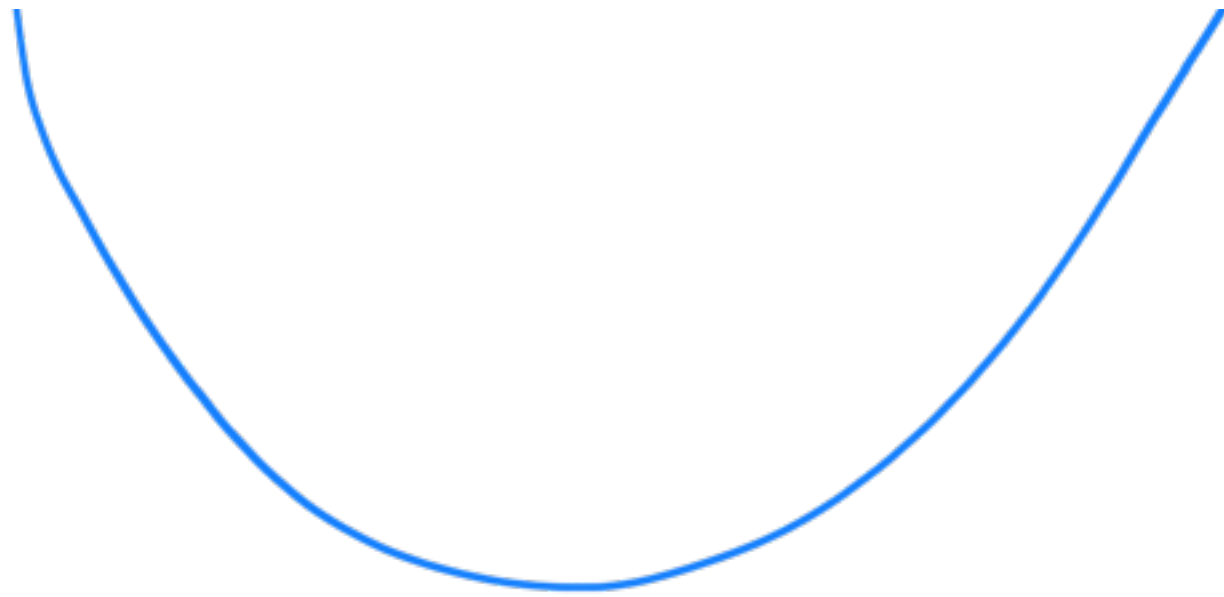
$$\Rightarrow -u'' + p u = \lambda u$$



$$u_m(x) = \left( \frac{2}{b-a} \right)^{1/2} \sin \left( m\pi \frac{x-a}{b-a} \right)$$

$$\lambda_m = E[u_m] = \frac{1}{2m} \left( \frac{m\pi}{b-a} \right)^2$$





$$-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} u + \frac{m\omega^2}{2} x^2 u = \lambda u$$

$$= \hbar\omega \left[ \frac{m\omega}{2\hbar} \left( x - \frac{\hbar}{m\omega} \frac{d}{dx} \right) \left( x + \frac{\hbar}{m\omega} \frac{d}{dx} \right) + \frac{1}{2} \right] u$$

ground state

$$\left( x + \frac{\hbar}{m\omega} \frac{d}{dx} \right) \langle x|0\rangle = 0$$

$$u_0(x) = \langle x|0\rangle = \left( \frac{m\omega}{\pi\hbar} \right)^{\frac{1}{4}} e^{-\frac{m\omega x^2}{2\hbar}}$$

$$u_n(x) = \frac{1}{\sqrt{n!}} \left( \frac{m\omega}{2\hbar} \right)^{\frac{n}{2}} \left( x - \frac{\hbar}{m\omega} \frac{d}{dx} \right)^n u_0(x).$$

$$a = \sqrt{\frac{m\omega}{2\hbar}} \left( x + \frac{i}{m\omega} p \right)$$

$$a^\dagger = \sqrt{\frac{m\omega}{2\hbar}} \left( x - \frac{i}{m\omega} p \right)$$

$$[a, a^\dagger] = 2 \frac{m\omega}{2\hbar} \frac{\hbar}{m\omega} = 1 \quad [q, p] = i\hbar$$

$$|m\rangle = \frac{1}{\sqrt{m!}} a^{\dagger m} |0\rangle \quad [x, p] = i\hbar$$

$$\psi_m(x) = \langle x | m \rangle.$$

$$a|z\rangle = z|z\rangle \\ z = x + iy$$

$$0 = x^2 y'' + x p(x) y' + q(x) y$$

where  $x p(x) = x^2 P(x)$

$$q(x) = x^2 Q(x)$$

and  $0 = y'' + P(x) y' + Q(x) y.$

$$y(x) = x^r \sum_{n=0}^{\infty} a_n x^n = \sum a_n x^{n+r}$$

$$y'(x) = \sum (n+r) a_n x^{n+r-1}$$

$$y''(x) = \sum (n+r)(n+r-1) a_n x^{n+r}$$

$$\text{So } 0 = x^2 y'' + x p y' + q y \quad \text{is}$$

$$0 = \sum (n+r)(n+r-1) a_n x^{n+r} + p(x) \sum (n+r) a_n x^{n+r} + q(x) \sum a_n x^{n+r}$$

$$\text{So } 0 = \sum \left[ (n+r)(n+r-1) + p(x)(n+r) + q(x) \right] a_n x^{n+r}$$

The lowest power of  $x$  here is for  $n=0$  given by the indicial equation:

$$0 = r(r-1) + p(0)r + q(0)$$

$$0 = r^2 + (p(0)-1)r + q(0)$$

$$r = \frac{1 - p(0) \pm \sqrt{(1 - p(0))^2 - 4q(0)}}{2}$$

Fuchs said use bigger root

$$r = \frac{1 - p(0) + \sqrt{(1 - p(0))^2 - 4q(0)}}{2}$$



But now if  $\infty$

$$p(x) = \sum_{j=0}^{\infty} p_j x^j$$

$$q(x) = \sum_{j=0}^{\infty} q_j x^j$$

So set  $m+n+j = k+r$   
 $m+j = k$   $m = k-j$

So  $j$  can't exceed  $k$ .

---

$$f(x) = \sum_{i=1}^{\infty} a_i u_i(x)$$

$$\int_a^b \rho u_j(x) f(x) dx = \int_a^b \sum_{i=1}^{\infty} a_i u_i(x) u_j(x) \rho dx$$

$$= \sum_{i=1}^{\infty} a_i \delta_{ij} = a_j$$

$$f(x) = \int_a^b f(y) \delta(x-y) dy$$

e.g.

$$P(\lambda) = 4\pi r$$



$$\Phi \quad \lambda$$
$$\lambda = \frac{h}{P} \quad \text{or} \quad \frac{h}{P}$$

$$P = \frac{h}{\lambda} \quad \text{or} \quad \frac{h}{\lambda}$$

$$p = \frac{1}{\lambda}$$

if we want to understand Nature

at  $\lambda \rightarrow 0$

$$p = \frac{1}{\lambda} \rightarrow \infty$$

$$N = \int_a^b dx \, u_k(x) \left( f(x) - \sum_{j=1}^n c_j u_j(x) \right)$$

$$N(r_n, u_k) = 0$$

$$E(r_n, r_n) = E(f, f)$$

bounded

$$-2 \sum c_k \lambda_k c_k$$

$$+ \sum C_{1k} C_{2k} \delta_{1k} \delta_{2k}$$

$$= \underbrace{E(f, f)}_{\text{norm}} - \sum \delta_{1k} C_{1k}^2$$

$$E(r_n, r_n) \leq E(f, f)$$

$$\|r_n\|^2 \leq \underbrace{E(f, f)}_{\text{bnd.}}$$

$n \rightarrow \infty$

$\rightarrow$  zero

$f \in D$

$$f(x) = \sum_{i=1}^n a_i(x) C_i$$

$$0 \leftarrow \int_a^b |f_n(x) - f(x)| p dx$$

$n \rightarrow \infty$

$$f(x) = \sum c_i u_i(x)$$

$$c_i = \int u_i(y) f(y) p(y) dy$$

$$f(x) = \sum u_i(x) \int u_i(y) f(y) p(y) dy$$

$$= \int f(y) \left[ \sum u_i(x) u_i(y) p(y) \right] dy$$

$$= \int f(y) \delta(x-y) dy$$

$$\delta(x-y) = \sum_{i=1}^{\infty} u_i(x) u_i(y)$$

$$\delta(x-y) = \sum_{i=1}^{\infty} u_i(x) u_i(y)$$

$$= \rho(x) \sum_{i=1}^{\infty} u_i(x) u_i(y)$$

$$\delta(x-y) = \rho(x) \rho(y) \sum_{i=1}^{\infty} u_i(x) u_i(y)$$

$$\mathcal{L} u = - (xu')' + \frac{a^2 u}{x} = \lambda x u$$

$$\lambda = -\zeta_{n,k}$$

$$\square = \Delta - \frac{\partial^2}{\partial t^2} = \Delta - \partial_0^2$$

$x^0 = ct$

$\vartheta(x)$



$$\theta(x) = \begin{cases} 0 & x < 0 \\ 1 & x > 0 \end{cases}$$

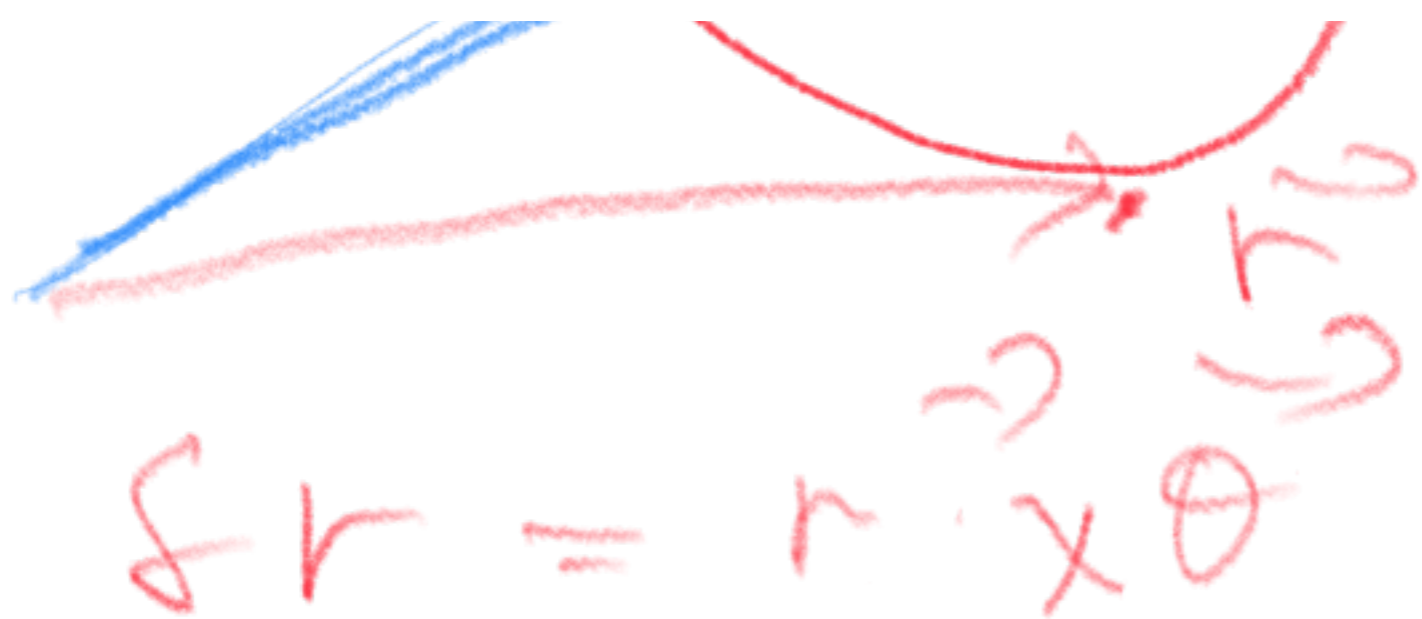
$$\partial_0 J^0 + \vec{\nabla} \cdot \vec{J} = 0$$

$$\partial_0 J^0 = -\vec{\nabla} \cdot \vec{J}$$

$$\equiv$$

$$\frac{\partial J^0}{\partial t}$$





$$-[(1-x^2)P_n']' = n(n+1)P_n$$

$$-(1-x^2)P_n'' + 2xP_n' = n(n+1)P_n$$

$$P_n'' = + \frac{2x}{(1-x^2)} P_n' - \frac{n(n+1)P_n}{(1-x^2)}$$

$$Q = \frac{n(n+1)}{(1-x^2)}$$

$$Q(x-1)^2 = \frac{(x-1)^2 n(n+1)}{(1-x)(1+x)}$$

$$= \frac{n(n+1)(x-1)}{1+x} \rightarrow 0 \quad x \rightarrow 1$$



$$z = \frac{1}{x}$$

$$\lim_{z \rightarrow 0} \frac{2z - \frac{z}{z}}{\frac{(-1/z^2)}{z^2}} = \frac{z}{z} - \frac{z}{z(z^2-1)} = 0$$

$$Q = \frac{1}{1-x^2} = \frac{1}{1-\frac{1}{z^2}}$$

$$a_{j+1} = \frac{j(j+1) - \lambda}{(j+2)(j+1)} a_{j-1}$$

$$a_{j+2} =$$

$$b_j = c_k^{-1} k_j$$

$$b = c T$$

$$b \psi^{-1} = \mathcal{L}$$

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$$\langle \chi | H | \psi \rangle = E_{\psi} \langle \chi | \psi \rangle$$

$$= \langle H \chi | \psi \rangle = E_{\chi} \langle \chi | \psi \rangle$$

$$E_{\psi} = E_{\chi}$$

unless  $\langle \chi | \psi \rangle = 0$ .

$$\mathcal{L} u = - (xu')' + \frac{m^2}{x} u = \lambda x u$$

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$$\mathcal{L} = \mathcal{L}(\phi_i, \partial_a \phi_i)$$

$$\frac{\partial \mathcal{L}}{\partial x^a} = 0$$

$$0 = \partial_a \left( \frac{\partial \mathcal{L}}{\partial \partial_a \phi} \delta \phi \right)$$

$$\nabla_i \phi \in \psi_k \theta_j \chi_k$$

$$= \epsilon_{jki} \nabla_i \phi \theta_j \chi_k$$

$$d\sigma^2 = \int_{ij=0}^3 g_{ij} dx^i dx^j$$

$$x^i = x^i(x^0, x^1, x^2, x^3)$$

i = 0, 1, 2, 3

2 - 3/11/21

