## 8

## Integral equations

### 8.1 Differential equations as integral equations

Differential equations when integrated become integral equations with builtin boundary conditions. Thus if we integrate the first-order ODE

$$
\begin{equation*}
\frac{d u(x)}{d x} \equiv u_{x}(x)=p(x) u(x)+q(x) \tag{8.1}
\end{equation*}
$$

then we get the integral equation

$$
\begin{equation*}
u(x)=\int_{a}^{x} p(y) u(y) d y+\int_{a}^{x} q(y) d y+u(a) . \tag{8.2}
\end{equation*}
$$

To transform a second-order differential equation into an integral equation, we use Cauchy's identity (exercise 8.1)

$$
\begin{equation*}
\int_{a}^{x} d z \int_{a}^{z} d y f(y)=\int_{a}^{x}(x-y) f(y) d y \tag{8.3}
\end{equation*}
$$

which is a special case of his formula for repeated integration

$$
\begin{equation*}
\int_{a}^{x} \int_{a}^{x_{1}} \cdots \int_{a}^{x_{n-1}} f\left(x_{n}\right) d x_{n} \ldots d x_{2} d x_{1}=\frac{1}{(n-1)!} \int_{a}^{x}(x-y)^{n-1} f(y) d y . \tag{8.4}
\end{equation*}
$$

We first write the second-order ODE in self-adjoint form $-\left(p u^{\prime}\right)^{\prime}+q u=\lambda \rho u$ as outlined in section 7.32 and then as $(p u)^{\prime \prime}=\left(p^{\prime} u\right)^{\prime}+(q-\lambda \rho) u$ which we integrate twice to

$$
\begin{align*}
p(x) u(x)= & p(a) u(a)+(x-a) p(a) u^{\prime}(a)+\int_{a}^{x} p^{\prime}(y) u(y) d y \\
& +\int_{a}^{x} d y \int_{a}^{y} d z(q(z)-\lambda \rho(z)) u(z) . \tag{8.5}
\end{align*}
$$

We then use Cauchy's identity (8.3) to integrate this equation to

$$
\begin{equation*}
p(x) u(x)=f(x)+\int_{a}^{x} k(x, y) u(y) d y \tag{8.6}
\end{equation*}
$$

in which $f(x)=p(a)\left[u(a)+(x-a) u^{\prime}(a)\right]$ and

$$
\begin{equation*}
k(x, y)=p^{\prime}(y)+(x-y)[q(y)-\lambda \rho(y)] . \tag{8.7}
\end{equation*}
$$

Example 8.1 (Legendre's equation) The function $p(x)=1-x^{2}$ in Legendre's equation $-\left[\left(1-x^{2}\right) P_{n}^{\prime}\right]^{\prime}=n(n+1) P_{n}$ vanishes at the end point $x=a=-1$ of the interval $[-1,1]$, so $f(x)$ also vanishes, and therefore formulas (8.6 and 8.7) give Legendre's integral equation as

$$
\begin{equation*}
\left(1-x^{2}\right) P_{n}(x)=-\int_{-1}^{x}[2 y+n(n+1)(x-y)] P_{n}(y) d y . \tag{8.8}
\end{equation*}
$$

Example 8.2 (Bessel's equation) The function $p(x)=x$ in Bessel's equation (10.11) $-\left[x J_{n}^{\prime}(x)\right]^{\prime}+\left(n^{2} / x\right) J_{n}(x)=x J_{n}(x)$ for $k=1$ vanishes at the end point $x=a=0$ of the interval $[0,1]$, so $f(x)$ also vanishes, and therefore since $q(x)=n^{2} / x$, formulas ( 8.6 and 8.7 ) give Bessel's integral equation as

$$
\begin{equation*}
x J_{n}(x)=\int_{0}^{x}\left\{1+(x-y)\left[n^{2} / y-y\right]\right\} J_{n}(y) d y \tag{8.9}
\end{equation*}
$$

In some physical problems, integral equations arise independently of differential equations. Whatever their origin, integral equations tend to have properties more suitable to mathematical analysis because derivatives are unbounded operators.

### 8.2 Fredholm integral equations

An equation of the form

$$
\begin{equation*}
\int_{a}^{b} k(x, y) u(y) d y=\lambda u(x)+f(x) \tag{8.10}
\end{equation*}
$$

for $a \leq x \leq b$ with a given kernel $k(x, y)$ and a specified function $f(x)$ is an inhomogeneous Fredholm equation of the second kind for the function $u(x)$ and the parameter $\lambda$. (Erik Ivar Fredholm, 1866-1927).

If $f(x)=0$, then it is a homogeneous Fredholm equation of the

