

$\exp(-st)$ lets many functions $F(t)$ that are not integrable over the half line $[0, \infty)$ have well-behaved Laplace transforms.

For instance, the function $F(t) = 1$ is not integrable over the half line $[0, \infty)$, but its Laplace transform

$$f(s) = \int_0^{\infty} dt e^{-st} F(t) = \int_0^{\infty} dt e^{-st} = \frac{1}{s} \quad (4.130)$$

is well defined for $\text{Re } s > 0$ and square integrable for $\text{Re } s > \epsilon$.

The function $F(t) = \exp(kt)$ diverges exponentially for $\text{Re } k > 0$, but its Laplace transform

$$f(s) = \int_0^{\infty} dt e^{-st} F(t) = \int_0^{\infty} dt e^{-(s-k)t} = \frac{1}{s-k} \quad (4.131)$$

is well defined for $\text{Re } s > k$ with a simple pole at $s = k$ (section 6.10) and is square integrable for $\text{Re } s > k + \epsilon$.

The Laplace transforms of $\cosh kt$ and $\sinh kt$ are

$$f(s) = \int_0^{\infty} dt e^{-st} \cosh kt = \frac{1}{2} \int_0^{\infty} dt e^{-st} (e^{kt} + e^{-kt}) = \frac{s}{s^2 - k^2} \quad (4.132)$$

and

$$f(s) = \int_0^{\infty} dt e^{-st} \sinh kt = \frac{1}{2} \int_0^{\infty} dt e^{-st} (e^{kt} - e^{-kt}) = \frac{k}{s^2 - k^2}. \quad (4.133)$$

The Laplace transform of $\cos \omega t$ is

$$f(s) = \int_0^{\infty} dt e^{-st} \cos \omega t = \frac{1}{2} \int_0^{\infty} dt e^{-st} (e^{i\omega t} + e^{-i\omega t}) = \frac{s}{s^2 + \omega^2} \quad (4.134)$$

and that of $\sin \omega t$ is

$$f(s) = \int_0^{\infty} dt e^{-st} \sin \omega t = \frac{1}{2i} \int_0^{\infty} dt e^{-st} (e^{i\omega t} - e^{-i\omega t}) = \frac{\omega}{s^2 + \omega^2}. \quad (4.135)$$

Example 4.13 (Lifetime of a Fluorophore) Fluorophores are molecules that emit visible light when excited by photons. The probability $P(t, t')$ that a fluorophore with a lifetime τ will emit a photon at time t if excited by a photon at time t' is

$$P(t, t') = \tau^{-1} e^{-(t-t')/\tau} \theta(t-t') \quad (4.136)$$

in which $\theta(t-t') = (t-t' + |t-t'|)/2|t-t'|$ is the Heaviside function. One way to measure the lifetime τ of a fluorophore is to modulate the exciting laser beam at a frequency $\nu = 2\pi\omega$ of the order of 60 MHz and to detect the phase-shift ϕ in the light $L(t)$ emitted by the fluorophore. That light is

the integral of $P(t, t')$ times the modulated beam $\sin \omega t$ or equivalently the convolution of $e^{-t/\tau} \theta(t)$ with $\sin \omega t$

$$\begin{aligned} L(t) &= \int_{-\infty}^{\infty} P(t, t') \sin(\omega t') dt' = \int_{-\infty}^{\infty} \tau^{-1} e^{-(t-t')/\tau} \theta(t-t') \sin(\omega t') dt' \\ &= \int_{-\infty}^t \tau^{-1} e^{-(t-t')/\tau} \sin(\omega t') dt'. \end{aligned} \quad (4.137)$$

Letting $u = t - t'$ and using the trigonometric formula

$$\sin(a - b) = \sin a \cos b - \cos a \sin b \quad (4.138)$$

we may relate this integral to the Laplace transforms of a sine (4.135) and a cosine (4.134)

$$\begin{aligned} L(t) &= -\tau^{-1} \int_0^{\infty} e^{-u/\tau} \sin \omega(u - t) du \\ &= -\tau^{-1} \int_0^{\infty} e^{-u/\tau} (\sin \omega u \cos \omega t - \cos \omega u \sin \omega t) du \\ &= \frac{\sin(\omega t) - \omega \tau \cos(\omega t)}{1 + (\omega \tau)^2}. \end{aligned} \quad (4.139)$$

Setting $\cos \phi = 1/\sqrt{1 + (\omega \tau)^2}$ and $\sin \phi = \omega \tau / \sqrt{1 + (\omega \tau)^2}$, we have

$$L(t) = \frac{1}{\sqrt{1 + (\omega \tau)^2}} (\sin \omega t \cos \phi - \cos \omega t \sin \phi) = \frac{1}{\sqrt{1 + (\omega \tau)^2}} \sin(\omega t - \phi). \quad (4.140)$$

The phase-shift ϕ then is given by

$$\phi = \arctan(\omega \tau) \leq \frac{\pi}{2} \quad (4.141)$$

and the lifetime of the fluorophore in terms of the phase-shift ϕ

$$\tau = (1/\omega) \tan \phi \quad (4.142)$$

which is much easier to measure than the lifetime τ . \square

4.10 Inversion of Laplace transforms

How do we invert the Laplace transform

$$f(s) = \int_0^{\infty} dt e^{-st} F(t)? \quad (4.143)$$

First we extend the Laplace transform from real s to $s + iu$

$$f(s + iu) = \int_0^{\infty} dt e^{-(s+iu)t} F(t) \quad (4.144)$$